

# Conductance and application of organic molecule pairs as nanofuses

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We propose that a pair of organic molecules can mimic the behavior of a macroscopic fuse at nanoscale, one component of the pair being the *on* state and the other the *off* state. For this task we make use of density-functional theory to calculate the physical properties of selected molecules, which have also been synthesized by our team. By this means we obtain the transmission spectra and the current of the proposed devices, which allows us to compare the behavior of the *on* and *off* states. Of particular interest is the *on/off* switch ratios, defined as the current ratios of the *on* and *off* structures at the corresponding bias voltage. In a first stage, we examine the best linker between the device and the electrode for high *on/off* switch ratios. Once this is determined, we test the influence of the electron richness of the system to provide a high *on/off* switch ratio. The entire analysis is also supported by the molecular projected self-consistent Hamiltonian, which provides a good way of understanding the molecular behavior. All the calculations support that interesting *on/off* switch ratios of two orders of magnitude could be obtained with these prototypical nanofuses.

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## I. INTRODUCTION

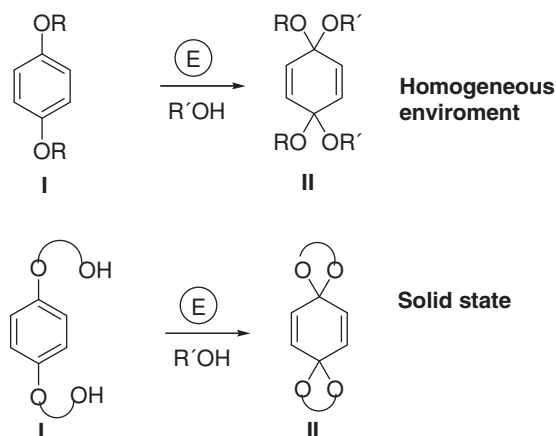
Electron transport through conjugated organic molecules is a central issue in molecular electronics and in material science on the nanoscale.<sup>1–14</sup> These molecules have critical properties for high conductance, such as small highest occupied molecular orbital (HOMO)–lowest unoccupied molecular orbital (LUMO) gaps and fully delocalized  $\pi$  systems. They can be used, for example, to connect electron reservoirs at the molecular level. In fact, in nature there are excellent examples of conjugated molecules facilitating the electronic connection between metallic centers in enzymatic systems.<sup>15–17</sup> Control of the flow of electrons between the two ends of molecular conductors is valuable in many scientific fields and may open up the possibility of simulating the behavior of standard circuitry at the molecular level.<sup>18,19</sup>

Some successful approaches to molecular switches based on metallic complexes have been reported using a difference of voltage to promote a change in the oxidation state of the metallic center. If we define the *on* state as the state of maximum conductance and the *off* state as the state of minimum conductance, they work in the *off/on* sense when the stimulus is increased, and consequently are voltage dependent. Moreover, they are reversible,<sup>20</sup> i.e., these switches change their conductivity when a specific voltage is applied but recover their initial features when the voltage is removed. The development of systems working in the *on/off* sense with an irreversible behavior to the voltage, mimicking a macroscopic fuse, is much less extensive. Ideally, when situated in a molecular circuit, these could (i) detect an overvoltage, thus protecting complex electronic systems, (ii) trigger other processes in a controlled manner, and (iii) act as binary components of read-only memories (ROMs).<sup>21</sup>

These factors prompted us to investigate this class of electronic device. In principle, a simple molecular fuse has to mimic the behavior of its macroscopic analogs.

Therefore we have imposed the following key features for an acceptable molecular fuse: (i) an easy integration into other previously reported “molecular wires,” (ii) to be able to drop its conductance at a desirable voltage with a high ratio between the conductivity of the *on* and *off* states, (iii) in some applications it should be desirable to be “switched” to the *on* state again in a controlled manner using some chemical or physical process, (iv) to be chemically stable in the *off* state at high voltage and, (v) to keep the same geometry and length in the *on* and *off* states in order to preserve the global structure of the system. This is an essential feature because it prevents any mechanical stress on the system when the switching process takes place. For example, photoswitches based on dithienylethene derivatives have the drawback that despite having excellent switching properties, they undergo changes in length and geometry during the switching event and a subsequent reorganization in the junction occurs to reach the final equilibrium state.<sup>22</sup>

With these requirements in mind we focused our attention on *p*-dialkoxybenzenes (I), which can be electrochemically oxidized to quinone bisketals (II). In this case, two oxygen atoms must be incorporated to the structure (Fig. 1).<sup>23,24</sup> The source of such oxygen atoms can be diverse. Thus, for example, in homogeneous environments oxygenated solvents or additives, such as methanol, are usually employed as an oxygen-atom source. Another solution is to incorporate the oxygen atom to the molecular framework at the expense of increasing the complexity of the global system. Nevertheless, such more complex “molecular fuses” are required for solid-state devices, owing to the absence of external sources of appropriate oxygen atoms. It is also worth noting that in the global reaction two protons are also lost into the homogeneous environment. In solid-state devices the fate of the hydrogen atoms is open to debate. It has been suggested that they

FIG. 1. Anodic oxidation of *p*-dialkoxybenzenes.

are lost as molecular hydrogen or as metal hydrides on the surfaces.<sup>25</sup>

A careful inspection of these two related compounds I and II shows that they fit with some of the key characteristics of a “molecular fuse” mentioned above and could be considered as candidates for the first approximation to a real nanofuse. First of all, type-I aromatic compounds have been extensively used in the construction of molecular electronic devices and therefore a trouble-free incorporation into elaborate molecular structures would be anticipated [requirement (i)].<sup>26–32</sup> Their  $\pi$ -conjugated systems can be also integrated in other extended  $\pi$ -conjugated molecular frameworks. On the other hand, the oxidation of the aromatic core of type-I structures provokes a disruption of the  $\pi$  conjugation and an increase in the HOMO-LUMO gap, hindering any potential electron flow through nonaromatic type-II structures. Interestingly, I to II conversion takes place at a fixed voltage, which in principle can be modulated in these simple systems changing the electron richness of the starting *on* type-I structures [requirement (ii)]. Furthermore, some procedures have been described to obtain conducting structures I from the corresponding insulating structures II [requirement (iii)].<sup>33–35</sup> Insulating quinone bisketals (II) are expected to be electrochemically inert within a voltage window of at least 4 V owing to their relatively stabilized HOMO and LUMO levels [requirement (iv)].<sup>36</sup> Finally, the basic structural cores of I and II are, respectively, a benzene and a tetrasubstituted 1,4-cyclohexadiene ring, which have essentially the same size and geometry [requirement (v)].<sup>37</sup> Another interesting property of our proposed system is that the starting and final *on* and *off* states are chemically stable and not transient states and could be synthesized and characterized independently and their corresponding properties measured. In fact, we have recently prepared some of these structures and proved experimentally their main hypothesized characteristics.<sup>38</sup>

Therefore, in principle, I/II pairs could be considered as a prototype of a molecular fuse. Nevertheless, some relevant issues should first be analyzed to prove their suitability. In particular, information about the *on/off* switch ratios, defined as the current ratios of the *on* and *off* structures at the corresponding bias voltage, is mandatory for this device. In this work we have studied this aspect computationally

focusing on (i) the role played by the anchoring groups between the molecule and the electrodes and (ii) the influence of the substituents on such *on/off* switch ratios, and bearing in mind that the presented results could be used to select the best starting structures I for subsequent experimental measurements of molecular conductivity.

The rest of this paper is organized as follows: In Sec. II we describe the molecular fuses analyzed and argue for the selection of this particular set of molecules. In Sec. III we explain the computational methodology applied to obtain the data. The results are shown in Sec. IV, in which qualitative information about the expected switching voltage is also extracted. Finally we draw our conclusions.

## II. MOLECULAR FUSES

In this work we have focused our attention on the capabilities for electron transport through molecular junctions of structures I and II. Although now they are experimentally available and their conductivity could be measured, such measures are usually difficult to carry out and a prior selection of the molecular targets is mandatory.<sup>39,40</sup>

Our working hypothesis is summarized in Fig. 2: The initial conductor system I can be electrochemically transformed into the insulator system II by the application of a suitable voltage in the presence of an appropriate oxygen source. Although for the *on/off* switch event development of charge in the molecule derived from hole-based electron transport (ET) is required, we have focused our study on the coherent ET mechanism, especially under low bias, which would be the usual working voltage for the device.

Theoretical approaches to molecular conduction based on density-functional theory (DFT) are now available<sup>41,42</sup> and can be used to discriminate between the structures with the best profiles and to understand the mechanism of ET through molecular conductors. The main limitation is that the experimental data and the theoretical predictions are not always in good agreement<sup>43,44</sup> due to multiple factors, such as geometric fluctuations in molecule-electrode interfaces,<sup>45</sup> structurally nonideal electrode surface interactions,<sup>46</sup> fast modifications of the internal geometry of the molecule,<sup>47</sup> solvent-molecule interactions,<sup>48,49</sup> or an underestimation of the band gap.<sup>50,51</sup> Nevertheless, we are mainly interested in the evaluation of the *on/off* switch ratios. Therefore we think that by maintaining the same computational parameters (functional, basis set, energy shift, Monkhorst-Pack grid, and pseudopotentials) in both systems the theoretical results may be comparable.

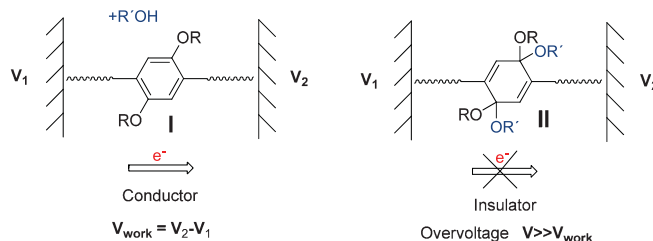


FIG. 2. (Color online) Working hypothesis.

From a practical point of view, the *on/off* switch ratio should be as high as possible. In the literature, different experimental switching ratios can be found: 2,<sup>52</sup> 25,<sup>53,54</sup> 100,<sup>55</sup> 250 (Ref. 56) (1 V),<sup>57</sup> 302,<sup>58</sup> 500 (0.55 V),<sup>59</sup> 1000,<sup>60</sup> or even examples in which the *off* state is literally “unplugged” to maximize the ratio.<sup>61</sup> However, the above described switching processes usually take place only from the *off* state to the *on* state when the external factor is enabled, and change to the *off* state again when that factor is removed. They are different from those studied in this paper because our proposed devices are able to work in the opposite *on/off* sense when increasing the voltage, and can also be maintained in the desirable state when the voltage is removed. To study the magnitude of *on/off* switch ratios, in this paper we have focused on the effects of two key parameters: (a) the nature of the linkers and (b) the effect of the substituents on the dialkoxybenzene moiety.

In Fig. 3(a) two different processes are shown to convert structure I to structure II. The one shown on the right needs the oxygen atom to be incorporated from an external source. However, for many solid-state devices no exogenous oxygen source is available and therefore more elaborate “molecular fuses” are required. The straightforward solution is the insertion into the structure of a pedant group including the required oxygen atom<sup>62</sup> as shown in the process on the left of this figure.

Based on these premises, we have selected some specific devices from those with structures shown in Fig. 3(b). They present *p*-terphenyl (TP) and expanded *p*-terphenyl structures

(ETP). These kinds of structures are highly valuable in fundamental studies of molecular conductivity,<sup>63,64</sup> their linear and rigid arrangements being highly appreciated in the field.

(E)TP-*on* are the *on* devices while (E)TP-*off* are their *off* states. The conductivity of the *on* states is mainly determined by the aromatic core of conducting structures I and a similar behavior is expected regardless of the substituents in the pedant alkoxy side chain. Therefore for simplicity we have also considered TP-*on* as the *on* states of the *off* devices TP-*off*2 when calculating the *on/off* switch ratios. Furthermore, it is known that the torsion angle plays a major role in the conductivity of the system.<sup>65</sup> Thus structures TP-*off*1 and TP-*off*2 can be used to determine the influence of the torsion angle on the conductivity. Similar reasoning has been used for the ETP-*on* devices and the *off* states ETP-*off*.

The presence of a group acting as linker between the electrodes and the molecule is indispensable in experimental electrode-molecule-electrode heterostructures.<sup>66</sup> The nature of such linker in contact with the metallic electrode and its ability to hybridize molecule and electrode electronic states determine the regime of electron transport.<sup>67</sup> For high conductivity, a strong-coupling regime is required. In this case, it has been described that linkers based on nitrogen,<sup>68–75</sup> sulphur,<sup>76–82</sup> phosphorus,<sup>83</sup> or carbon<sup>84</sup> can be used to carry out measurements of conductivity values.

To carry out the study of the linkers, we selected TP structures to calculate the *on/off* switch ratios of our proposed molecular fuses. Thus we included linkers with different electronic properties: (i) a sulphur-based electron-rich anchoring group (TPS), and (ii) two electron-poor nitrogen-based anchoring groups with different affinities to metals, such as nitrile (weak) TPCN and isonitrile (strong) TPNC.<sup>85,86</sup> As mentioned above, a problem in the study of TP structures is that they are nonplanar and the corresponding torsion angle between the aromatic rings is a key parameter in the conduction through such molecules. We have taken this effect into account when studying the conductivity through the conformations of minimum energy determined by our DFT calculations. This is a crucial point for us since we are also interested in building up the corresponding devices to validate our predictions and thus we must assume the most realistic situation in the calculation.

Once the linker analysis was made, in order to study the effect of substituents in the core aromatic ring, we selected ETP structures, which are almost planar in their ground states. The spacers between the aromatic rings ensure that the presence of the substituents does not significantly affect the torsion angle of the device and such effects can be almost eliminated.

### III. COMPUTATIONAL METHODOLOGY

One common approach in this kind of study is to build a conceptual device in which molecules are anchored between metallic electrodes in a sandwich-type junction. This configuration simulates the frequently used approach of scanning probe microscopy to measure molecular conductivity. In this area, there is abundant literature about systems using gold for metallic electrodes in both theoretical and experimental studies. For simplicity and to facilitate an experimental verification of our predictions, we selected bulk gold with

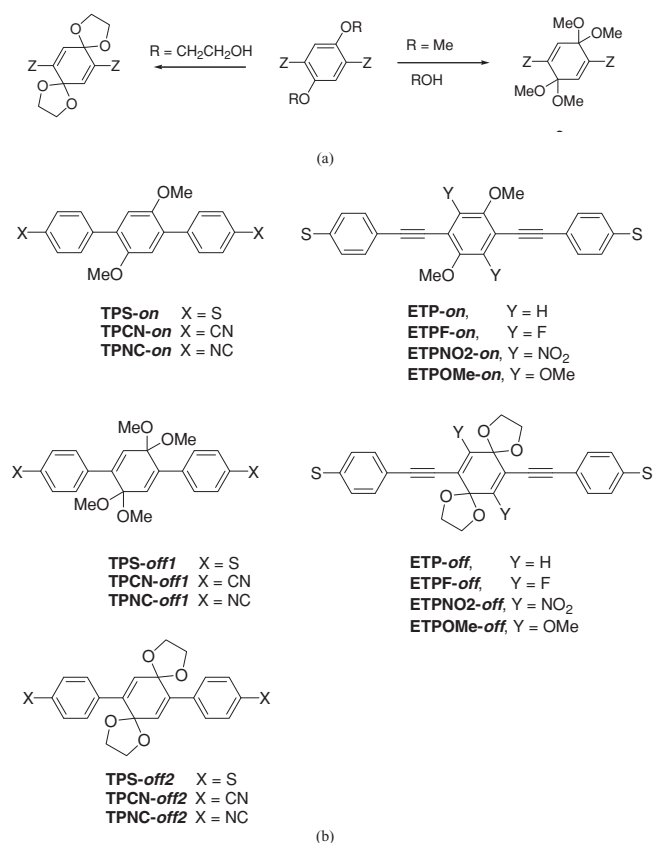


FIG. 3. (a) Processes to convert structure I to structure II. (b) Definition of the structures to be simulated.

Au(111) surfaces as electrodes throughout our study. Instead of simulating the Au electrodes with thousands of atoms, we simulated the Au electrode surfaces using a (4×4) cell with periodic boundary conditions. From this, we defined a supercell consisting of two layers of 32 and 48 atoms, respectively, on both sides of the scattering region. The size of the supercell ensures that there is no significant interaction between the molecules of the next neighboring supercells.

Gas phase TP and ETP structures were previously optimized by Gaussian 03 at the DFT level, using the B3LYP hybrid functional and the 6-31G(*d,p*) basis set for all atoms. They were then joined to these Au electrodes. However, the nature of the junction depends on the nature of the linker. The sulphur-Au junction has been extensively studied. Although the nature of the gold-sulphur bond is not clear, some recent advances have been achieved,<sup>87,88</sup> and many theoretical studies predict that “triangle” (interaction with two Au atoms) and “pyramid” (interaction with three Au atoms) interactions are the most favorable in terms of energy though small differences are usually obtained.<sup>89,90</sup> As our main purpose is to determine the switch of conductivity and not the real nature of Au-S interaction, we selected for simplicity the pyramid configuration. Although tilted structures are possible, we only considered configurations in which the molecule was placed perpendicular to the electrode surface above a threefold face-centered-cubic (fcc) hollow site, i.e., those above a threefold site beneath which there is no Au atom in the second layer but with an Au atom in the third layer. An additional gold layer was included in one of the electrodes to ensure that the two S-Au connections were in similar fcc sites. Moreover, in all the devices, we fixed the orientation of the central ring to intersect the Au face-centered-cubic triangular faces. This assumption simplifies the problem because it is known that in these systems the conductivity depends on the intersection of the plane of the molecule in the Au face-centered-cubic triangular faces.<sup>91</sup>

We used molecule TPS-*on* to find the optimal S-Au distance in the pyramidal configuration for all TPS-based devices. To this end, we fully relaxed the supercell, gradually increasing the electrode distance until forces were less than 0.04 eV/Å, using the Spanish Initiative for Electronic Simulations with Thousands of Atoms (SIESTA) code.<sup>92</sup> The geometry of the supercell was optimized at DFT level with Perdew-Burke-Ernzerhof generalized-gradient approximation (GGA-PEB) functional with a double- $\zeta$  plus polarization basis set for the organic molecule and gold atoms, with a 5-mRy energy-shift parameter, and a 5×5 Monkhorst-Pack grid for electronic structure calculations. During these calculations, core electrons and nuclei are replaced by Troullier-Martin’s norm-conservative pseudopotentials.<sup>93</sup> For the sake of simplicity, we selected an initial Au-Au distance of 2.88 Å in bulk Au and kept the Au layers fixed during the optimization process. This is a reasonable assumption due to the low distortion of the Au arrangement expected.<sup>94–96</sup> While SIESTA supports the spin-polarized treatment of electrons, taking into account the nature of our system we employed only unpolarized calculations in this study. With these initial conditions we found that a distance from sulphur to the Au contact site of 2.25 Å, corresponding to an S-Au distance of 2.82 Å, presents a minimum energy. This value is close to other values described

in similar Au-S arrangements.<sup>97–99</sup> For ETP-*on*, the minimum energy was obtained with a slightly shorter S-Au bond length of 2.68 Å.

Once the S-Au distance is determined, we know the electrode distance for the *on* devices, this being 19.86 Å for TPS-*on* and 24.52 Å for ETP-*on*. For the rest of the sulphur-containing devices (*off* states TP-*off*, *on* states ETPF-*on*, ETPNO2-*on*, ETPOMe-*on*, and all *off* states ETP), in order to mimic a real device with a fixed length, we considered that their electrode separations were also those corresponding to TPS-*on* and ETP-*on*. All were fully relaxed, yielding almost symmetric devices. The only exceptions were TPS-*off* devices, which presented two different distances from sulphur to the contact site (2.15 and 2.30 Å). It is worth noting that the linearity was almost conserved (TPS-*off*1, 176.10°; TPCN-*off*1, 177.50°).

The nitrile and isonitrile functionalized devices TPCN-*on* and TPNC-*on* were set up following similar reasoning. In this case, we placed the nitrogenated linker group in the *top* configuration. The optimal contact distances for TPCN-*on* and TPNC-*on* were 2.25 and 2.12 Å, respectively. Having established these distances, we set the rest of the devices.

With the optimized devices TP and ETP we carried out computational studies combining DFT and the nonequilibrium Green’s-function scattering approach provided by a Quantum Wise software package<sup>100</sup> to calculate the electron transport properties of our proposed molecular fuses. We used the GGA-PEB functional with the double- $\zeta$  plus polarization basis set for the organic molecule and the single- $\zeta$  plus polarization basis set for the Au electrodes. A 3×3 Monkhorst-Pack grid for electronic structure calculations and a 6×6 K sampling in the transmission spectra with a 5-mRy energy-shift parameter were selected, based on convergence criteria and computational cost. The values of these three parameters are well tested in similar studies.<sup>101</sup>

Current (*I*) was calculated as a function of the voltage using the Landauer formula,

$$I = \frac{2q}{h} \int_0^\infty T(E) [f(E - \mu_L) - f(E - \mu_R)] dE, \quad (1)$$

where  $T(E)$  is the transmission through the system as a function of the energy,  $f$  is the Fermi-Dirac distribution function,  $q$  is the charge of the electron,  $h$  is Planck’s constant, and  $\mu_{L(R)}$  are the Fermi level at the applied voltage at the left (right) contact. When an external voltage  $V$  is applied, the right voltage is  $V/2$  and the left contact voltage is set to  $-V/2$ .

## IV. RESULTS AND DISCUSSION

### A. Influence of the anchoring groups

In Fig. 4 we show a comparison between the transmission spectra of the three TP devices at low bias ( $\Delta V \approx 0$ ). The transport properties of the devices at low (operational) voltage ( $V$ ) are related to the nature of the transmission function at  $E_f \pm V/2$ ,  $E_f$  being the Fermi level of gold. Thus we are interested mainly in the proximities of such a level, taking into account that molecular electronic devices are expected to work under a relatively low voltage (less than 1 V).<sup>102</sup> However, for the sake of completeness, we show a wide voltage window of 4 eV around the Fermi level in the figure.



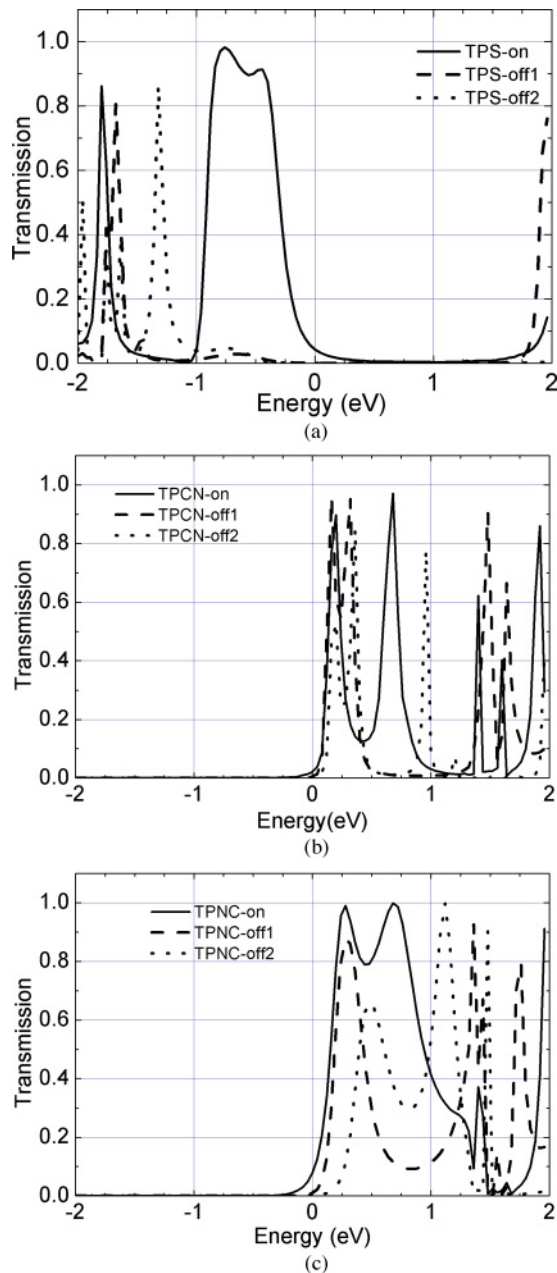


FIG. 4. (Color online) Transmission coefficients of TP devices at low bias.

When sulphur is used as linker, we can observe broad peaks, which correspond to a strong-coupling regime [Fig. 4(a)]. The main characteristic is that the transmission coefficients near the Fermi energy corresponding to the *on* device TPS-*on* are higher than those corresponding to *off* TPS-*off1* and TPS-*off2* devices. This is related to the presence of a broad, almost transparent resonance from  $-0.28$  to  $-0.94$  eV for the TPS-*on* device. An eigenchannel analysis showed that only one channel is open throughout the energy range studied ( $E_f \pm 2.0$  eV) in the three devices. Molecular projected self-consistent Hamiltonian (MPSH) analysis is a useful tool to clarify the role of the molecule in the transport behavior of the global system, relating or not to the transmission peaks with the molecule. Such analysis is carried out projecting the self-consistent

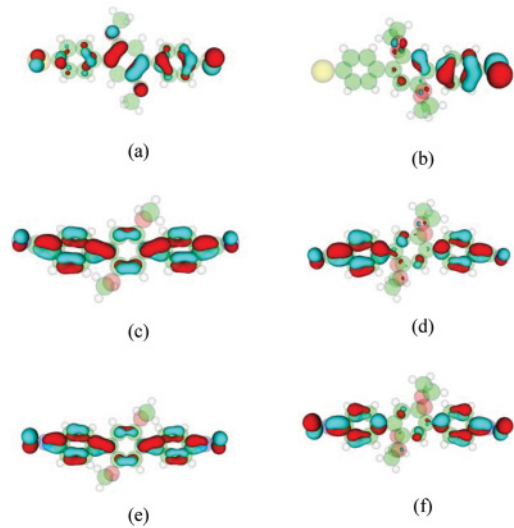


FIG. 5. (Color online) Isosurface plots of the MPSH wave functions for TPS devices at low bias ( $V = 0$ ) at different energies: (a) HOMO level of TPS-*on* device at  $-0.42$  eV; (b) HOMO level of TPS-*off2* device at  $-0.67$  eV; (c) LUMO level of TPCN-*on* device at  $0.14$  eV; (d) LUMO level of TPCN-*off2* device at  $0.14$  eV; (e) LUMO level of TPNC-*on* device at  $0.11$  eV; (f) LUMO level of TPNC-*off2* device at  $0.24$  eV.

Hamiltonian onto the atoms in the central region, and then diagonalized to produce an energy spectrum. This can be used to show how the molecular levels for a molecule are affected when placed between the two electrodes.<sup>100</sup>

Consequently with the above-mentioned results, a MPSH analysis of the TPS-*on* device showed that this broad transmission peak is derived from the presence of the HOMO level at  $-0.42$  eV (see Fig. 5) and the HOMO-1 level at  $-0.89$  eV. Moreover, we observed a high delocalized wave function of HOMO and HOMO-1 levels in the MPSH isosurface plots for the TPS-*on* device. Therefore a good connection between the metallic electrodes is achieved, maximizing the transmission coefficients.

Predicting the *on/off* switch voltage in solid state is not a simple task because of the broadening of the molecular levels due to the interaction with the electrodes and the alignment of such levels toward the Fermi energy of the electrode. In a simple approximation, the oxidation reaction for these substrates is expected to occur when the difference of voltage centered in the Fermi energy of the electrode includes the broadened HOMO level, being the maximum probability of such an *on/off* switch at the center of the broadened HOMO level. Therefore an estimated oxidation potential (approximately  $0.5$ – $0.9$  V) is predicted for such structures with sulphur/gold junctions. However, DFT usually overestimates HOMO values and higher oxidation potentials would be expected in real experiments.<sup>103</sup> This *on/off* voltage value is important because it limits the working voltage, which must always be smaller in order for the device to work properly. All these assumptions are limited by the fact that in this work we are considering only coherent transport in the final *on* and *off* states and not the inherent dynamic process of the switching event, which would be related to hopping mechanisms. In that

case, instant charges would be developed in the molecule and probably the transmission spectrum would change. Although the precise mechanism of *on/off* transformation is not known even in homogeneous environments, previous calculations suggest that the global geometry of the system is retained during the switching event despite the generation of charge in the molecule.<sup>38</sup>

In TPS-*off*1 and TPS-*off*2 devices the electron transport is also dominated by broad peaks (TPS-*off*1,  $-0.52$  to  $-0.82$  eV; TPS-*off*2,  $-0.52$  to  $-1.0$  eV) but with much lower transmission coefficients. In this case, the low transmission values suggest an unfavorable mixing between electronic levels of the metallic electrodes at these energies. A similar MPSH analysis showed that the molecule has energy levels at these energies, but they are highly localized in only half of it, as we initially proposed, based on the insulating character of type-II structures.<sup>104</sup> Although the HOMO and HOMO-1 levels should be degenerated in TPS-*off*1 and TPS-*off*2 devices, we commented previously that they are asymmetrically joined to the electrodes and this asymmetry is revealed in the corresponding alignment of such levels. This geometrical asymmetry has no significant consequences in the transport mechanism but will provoke asymmetric *I/V* curves.

It could be argued that the disconnection derives from differences in torsion angles in the three devices but this is not the sole reason: Device TPS-*off*1 ( $16.0^\circ$ ) presents a lower torsion angle than TPS-*on* ( $25.9^\circ$ ). Nevertheless, if we compare *off* devices TPS-*off*1 and TPS-*off*2, the expected relationship between the torsion angle and the conductivity is recovered (TPS-*off*1,  $16.0^\circ$ , TPS-*off*2,  $42^\circ$ ).

The use of nitrogenated linkers changed the profile of the transmission spectra. In these cases, the conduction is LUMO controlled.<sup>95,96</sup> We can also observe that sharper peaks are obtained in the case of nitrile ( $-\text{CN}$ ) containing devices TPCN because a weaker interaction between such groups and gold surfaces is expected. As in the case of TPS, only one eigenchannel is open in the voltage window studied. Moreover, the maximum of the transmission spectra is related to the presence of energy levels with a high contribution from the original organic molecule (TPCN-*on*, 0.14, 0.63 eV; TPCN-*off*1, 0.12, 0.28 eV; TPCN-*off*2, 0.14, 0.31 eV; TPNC-*on*, 0.11, 0.52 eV; TPNC-*off*1, 0.10, 0.22 eV; TPNC-*off*2, 0.24, 0.34 eV). The main difference from TPS devices is the relatively high transmission coefficients corresponding to the *off* structures. The MPSH isosurface plots for TPS, TPCN, and TPNC devices showed the reason for this behavior. TPCN and TPNC LUMO levels are delocalized throughout the entire molecule, because the decoupling of the two subunits in the *off* states is inefficient. Another difference compared with the TPS-*on* device is the expected *on/off* switch voltage, which is related to the HOMO level. In this case, such levels are located at  $-2.61$  eV for the TPCN-*on* structure and  $-2.56$  eV for the TPNC-*on* structure. In this situation a undesired *on/off* switch voltage of up to 6 V would be required to trigger the switching event.

Current can be also calculated. As we commented earlier, we are mainly interested in the low bias region because real operating systems will work under relatively low voltage. A voltage window of 1.5 V was therefore selected, taking into account also the experimental value of the oxidation peaks of terphenyls in solution.<sup>38</sup> The currents calculated at different

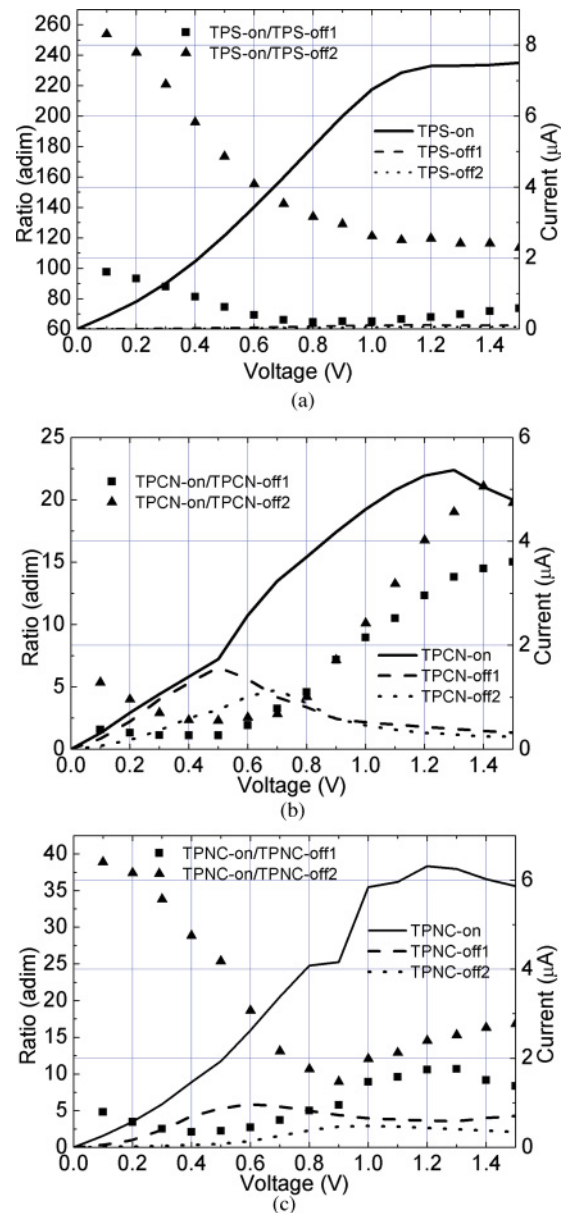


FIG. 6. (Color online) Current-voltage and switching values as voltage function.

voltages for TP devices are shown in the corresponding *I-V* curves of Fig. 6.

According to the transmission spectra, the conductivity of conducting TP-*on* structures is higher than that of insulating TP-*off* structures. This fact is especially significant for sulphur bonded structures. Nevertheless, devices with similar length and structure, such as TP-*off*1 and TP-*off*2, showed significant differences in conductivity, probably due to the different torsion angle values. This feature is worth noting because the torsion angle seems to be controlled by the nature of the side chains. Thus molecular fuses containing other more elaborated side chains might present a double switching process: a disconnection in the  $\pi$ -conjugated system and a higher torsion angle.

The *I/V* curves are highly symmetric with the exception of the *off* structures of the TPS-*on* device, owing to the previously

mentioned slight asymmetry in the anchoring groups. Even in this limit the differences in current are less than 7% and for the sake of clarity we have only shown the  $I/V$  curve in the forward direction in all cases. We can observe high  $dI/dV$  values at voltages lower than 1 V, which belong to electrons with energies corresponding to the maxima of the transmission functions. Negative differential conductance (NDC) was also observed in TPCN and TPNC devices. This fact is assigned to misalignments between the Fermi level of the electrodes and the molecular levels derived from the presence of voltage.

We have estimated the *on/off* switch ratios from 97 at 0.1 V to 71 at 1.5 V for the TPS-*on*/TPS-*off*1 devices, and 253 at 0.1 V to 116 at 1.5 V for the TPS-*on*/TPS-*off*2 devices. Lower values were obtained for the TPCN and TPNC devices. Thus, for example, a maximum value of 21 at 1.4 V was obtained for the TPCN-*on*/TPCN-*off*2 pair and a maximum value of 38 at 0.1 V for the TPNC-*on*/TPNC-*off*2 pair. The TP-*off*2 structures always show the best switching values, taking into account that they present higher torsion angles.

### B. Influence of the substituents

The substituents of ETP structures were selected to change the electronic richness of the key aromatic subunit.<sup>105,106</sup> This may have two consequences: (i) a change in the oxidation potential of the *on* state and (i) a change in the alignment of the molecular energy levels with the Fermi level of gold, which is essential for electron transport. The first is interesting, since it allows the selection of molecular fuses with different *on/off* switch voltages depending on the substituents. The second is obviously related to the efficiency of the device. Taking into account that the better switching profile was obtained with a sulphur atom as a linker, the study focused on structures containing such a linker. Concerning the substituents, in addition to model systems ETP-*on* and ETP-*off*, we selected devices with two common deactivating (i.e., electron-withdrawing) groups such as fluorine (ETPF-*on* and ETPF-*off*,  $R = F$ ) and nitro (ETPNO2-*on* and ETPNO2-*off*,  $R = -NO_2$ ), and one prototypical activating (i.e., electron releasing) group such as the methoxy group (ETPOMe-*on* and ETPOMe-*off*,  $R = OMe$ ).<sup>107</sup> Although fluorine is the least deactivating group among the halide family, it presents the least atomic volume and, consequently, the least distortion in the parent structure.

The transmission spectra of the ETP *on* devices, shown in Fig. 7, have broad, almost transparent resonances from  $-0.3$  to  $-1.0$  eV, similar to the previously described TPS-*on* device. The first maximum is again correlated with the presence of a MPSH level (ETP-*on*,  $-0.46$  eV; ETPF-*on*,  $-0.52$  eV; ETPNO2-*on*,  $-0.70$ ; ETPOMe-*on*,  $-0.40$  eV) (Fig. 8). The differences in energy of such MPSHs between ETP-*on* and the *modified*-ETP-*on* structures should depend on the activating and deactivating effects of the substituents in the aromatic ring. Nevertheless, if we consider the ETP-*on* device as a reference, we could only observe a significant decrease in the energy of the HOMO level of the ETPNO2-*on* device (0.24 meV). The other variations of 0.06 meV are well within error limits in the DFT calculations and cannot be certainly attributed to the electronic effects of the substituents.

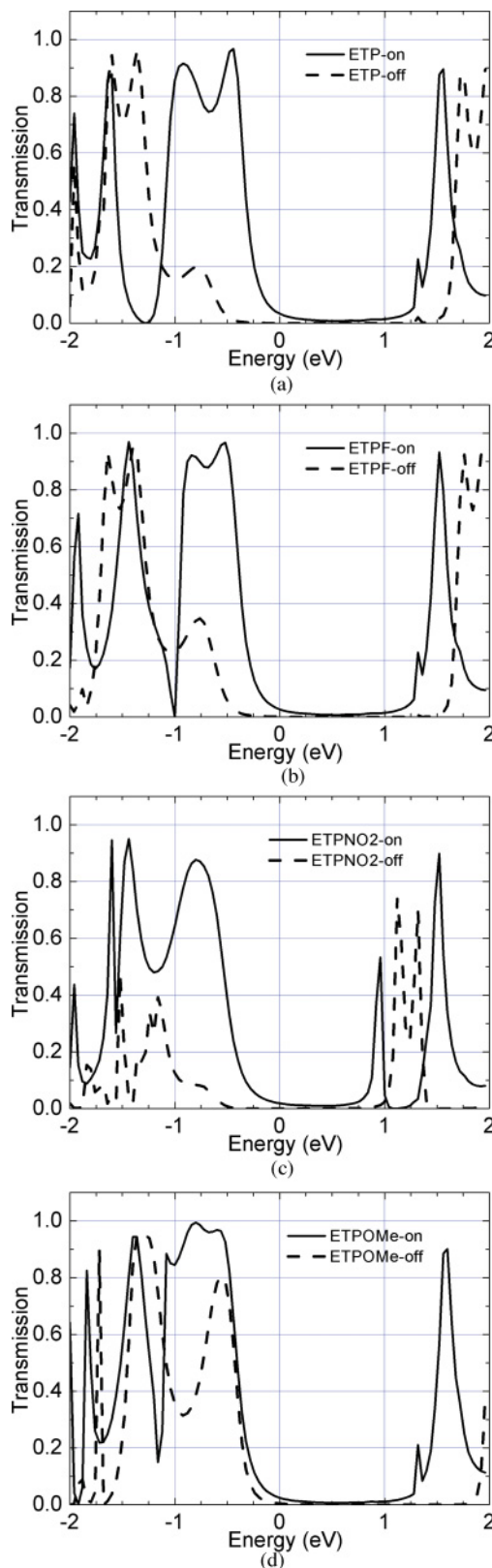


FIG. 7. (Color online) Transmission spectra of ETP devices at low bias.

Concerning the *off* devices, the lowest transmission coefficients at low bias were obtained with the ETPNO2-*off* device, which includes the highest deactivating group



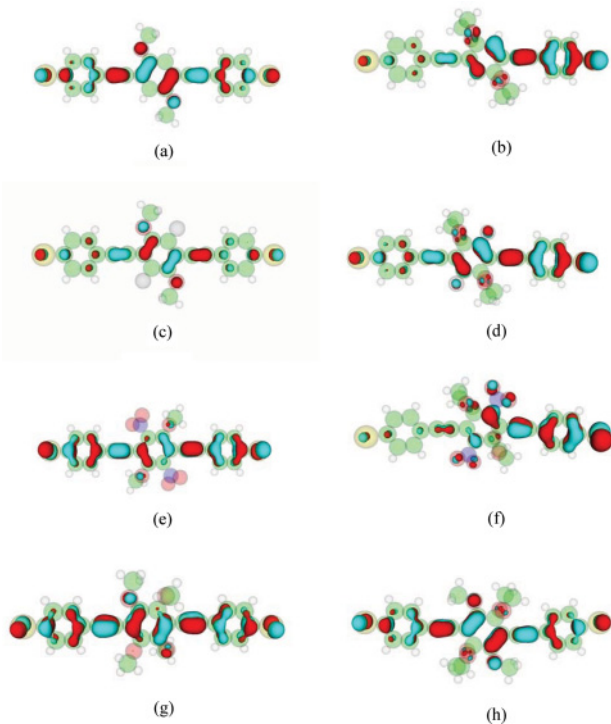


FIG. 8. (Color online) Isosurface plots of the MPSH wave functions for ETP devices at low bias ( $V = 0$ ) at different energies: (a) HOMO level of ETP-*on* device at  $-0.46$  eV; (b) HOMO level of ETP-*off* device at  $-0.76$  eV; (c) HOMO level of ETPF-*on* device at  $-0.52$  eV; (d) HOMO level of ETPF-*off* device at  $-0.74$  eV; (e) HOMO level of ETPNO2-*on* device at  $-0.70$  eV; (f) HOMO level of ETPNO2-*off* device at  $-0.71$  eV; (g) HOMO level of ETPOMe-*on* device at  $-0.40$  eV; (h) HOMO level of ETPOMe-*off* device at  $-0.49$  eV.

studied. Consequently, the highest transmission coefficients correspond to the ETPOMe-*off* device, with the activating  $-OMe$  group. The fluorine-containing device ETPF-*off* and model device ETP-*off* showed intermediate values, with that corresponding to ETPF-*off* being slightly higher. An increase of the electron density in the *off* structures seems to derive an inefficient disconnection between the two subunits. The MPSH analysis of ETP devices demonstrates this fact. Two MPSH levels are present within the voltage window in all *off* ETP devices (ETP-*off*,  $-0.76$  and  $-0.92$  eV; ETPF-*off*,  $-0.74$  and  $-0.94$  eV; ETPNO2-*off*,  $-0.71$  and  $-0.81$  eV; ETPOMe-*off*,  $-0.49$  and  $-0.73$  eV). The nondegeneracy of HOMO and HOMO-1 levels in *off* ETP devices again reveals small asymmetries in the geometry of the device. Significant differences can be appreciated in the isosurface plots of the corresponding MPSH. The *on* ETP structures present a fully conjugated orbital while *off* structures are efficiently disconnected only in nitro- (ETPNO2-*off*) and hydrogen- (ETP-*off*) substituted devices.

The  $I$ - $V$  curves [Fig. 9(a)] and switching ratios [Fig. 9(b)] at different voltages were calculated for ETP-*on* and ETP-*off* devices. From a practical point of view, the ETPNO2-*on*/ETPNO2-*off* pair presents the best switching profile at low bias with *on/off* switch ratios ranging from 636 at 0.1 V to 52 at 1.5 V. Interestingly, useful ratios can also be obtained using

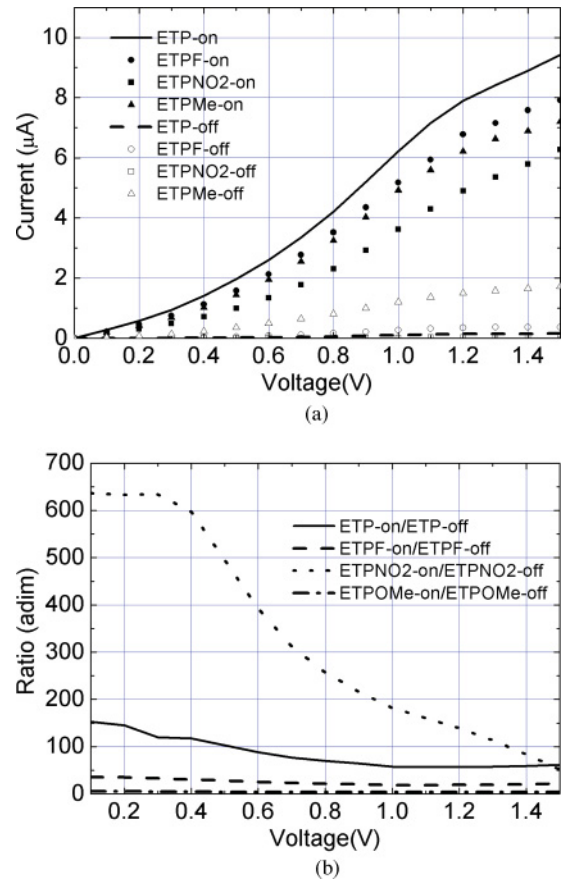


FIG. 9. (Color online) Current-voltage curves and switching values of ETP pairs as voltage function.

the simple unsubstituted ETP-*on*/ETP-*off* pair (152 at 0.1 V and 61 at 1.5 V).

## V. CONCLUSIONS

In summary, we have presented a study of the potential application of dialkoxybenzenes/quinone bisketal pairs as molecular fuses owing to their unique features, which could be of interest in the field of nanoelectronics. It is also worth noting that closely related structures have been synthesized recently, confirming some of the key predicted properties. Besides their remarkable geometrical behaviors we described in this paper interesting *on/off* switch ratios using in the modeling of the system SIESTA and Quantum Wise codes. The best *on/off* switch ratios were obtained when sulphur atoms were used as linkers between the molecule and the electrodes (TPS and ETP). In these particular cases the transmission spectra correspond to a system in which the molecular conductivity is HOMO controlled. Interestingly, in this regimen the *on/off* switch event could take place at a convenient voltage. With respect to the effect of the electron richness on the *on/off* switch ratio we found that strongly deactivating groups ( $-NO_2$ ) highly improve it. Thus, for example, the *on/off* switch ratio ranged from 636 for deactivated device ETPNO2, being only 7 for activated ETPOMe for the same voltage (0.1 V). As the *on/off* switch voltage is controlled by the position of the HOMO



level and it does not substantially change by the presence of different substituents, we predict a relatively small influence of such substituents on the *on/off* switch voltage. Therefore we can conclude that molecules presenting sulfur-based groups as linkers and electron withdrawing groups in the central aromatic ring would be interesting candidates as prototypical nanofuses for experimental measures. Although the present calculations suggest the viability of our proposal, a real device should increment the *on/off* switch ratio from the calculated two orders of magnitude to at least six orders of magnitude. Based on these precedents new calculations of a second generation of nanofuses are underway.

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