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Bracket map for Heisenberg group and the characterization of cyclic subspaces

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Abstract

The bracket map was originally considered in [7] for locally compact abelian groups. In this work we extend the study of bracket maps to the noncommutative setting, providing characterizations of bases and frames for cyclic subspaces of the Heisenberg group. We also indicate how to generalize these results to a class of non-abelian nilpotent Lie groups whose irreducible representations are square integrable modulo the center.¹

1 Introduction

Principal shift invariant subspaces obtained as the closed linear span of integer translates of a single function $\psi \in L^2(\mathbb{R}^d)$ have been studied recently in connection with the theory of Multiresolution Analysis for wavelets.

The general situation could be described as follows. Let \mathcal{G} be a topological group, and denote by $g \mapsto T_g$ a unitary representation of \mathcal{G} on a Hilbert space \mathcal{H} . The Principal Shift Invariant (PSI) subspace generated by $\psi \in \mathcal{H}$, denoted by $\langle \psi \rangle_{\mathcal{G}, T}$, is the closure in \mathcal{H} of the linear span generated by $\{T_g \psi\}_{g \in \mathcal{G}}$, and one could ask to characterize those $\psi \in \mathcal{H}$ so that the collection $\{T_g \psi\}_{g \in \mathcal{G}}$ is an orthonormal basis, a Riesz basis or a frame for $\langle \psi \rangle_{\mathcal{G}, T}$.

Such characterizations have been provided for countable abelian groups \mathcal{G} equipped with the discrete topology, whose unitary representations are dual integrable (see [7] for details). Such characterizations are given in terms of a bracket map

$$[\cdot, \cdot] : \mathcal{H} \times \mathcal{H} \rightarrow L^1(\hat{\mathcal{G}}, d\alpha)$$

where $\hat{\mathcal{G}}$ is the character group of \mathcal{G} with Haar measure $d\alpha$, such that

$$\langle \varphi, T_g \psi \rangle = \int_{\hat{\mathcal{G}}} [\varphi, \psi](\alpha) \overline{\alpha(g)} d\alpha \quad \forall \varphi, \psi \in \mathcal{H}. \quad (1.1)$$

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Two classical examples of bracket maps are the following. The first one is the map $k \mapsto T_k$ given by $T_k f(x) = f(x + k)$ as a unitary representation of $(\mathbb{Z}^d, +)$ on $L^2(\mathbb{R}^d)$. In this case

$$[\varphi, \psi](\xi) = \sum_{l \in \mathbb{Z}^d} \hat{\varphi}(\xi + l) \overline{\hat{\psi}(\xi + l)}, \quad \xi \in \mathbb{R}^d, \quad \varphi, \psi \in L^2(\mathbb{R}^d) \quad (1.2)$$

where $\hat{\psi}$ stands for the Fourier transform of ψ . The bracket map for this situation has appeared in [4, 5]. The second one is the Gabor unitary representation of the group $(\mathbb{Z}^d \times \mathbb{Z}^d, +)$ on $L^2(\mathbb{R}^d)$ given by

$$(k, l) \mapsto M_l T_k \psi(x) = e^{-2\pi i l \cdot x} \psi(x - k).$$

In this case

$$[\varphi, \psi](x, \xi) = Z\varphi(x, \xi) \overline{Z\psi(x, \xi)}, \quad x, \xi \in \mathbb{R}^d, \quad \varphi, \psi \in L^2(\mathbb{R}^d) \quad (1.3)$$

where

$$Z\psi(x, \xi) = \sum_{k \in \mathbb{Z}^d} e^{-2\pi i k \cdot \xi} \psi(x - k)$$

is the Zak transform of the function $\psi \in L^2(\mathbb{R}^d)$. The bracket map for this situation has appeared in [9, 8].

An appearance of the bracket map in disguise is in [1] where invariant subspaces for local fields are studied with respect to suitable defined translations.

As far as we know, no attempt has been done to provide a bracket map for non-Abelian groups. Our purpose in this paper is to provide such a map for left translations on the Heisenberg group. A characterization of frames and Riesz bases for left translations on nilpotent Lie groups whose irreducible representations are square integrable modulo the center (which include the Heisenberg group), also called SI/Z Lie groups, has been recently announced in [3]. These characterizations are in terms of range functions, rather than bracket maps.

In this paper the Heisenberg group \mathbb{H}^d is identified with $\mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}$ and noncommutative group law given by

$$(p, q, t)(p', q', t') = (p + p', q + q', t + t' + p \cdot q'). \quad (1.4)$$

where \cdot stands for \mathbb{R}^d scalar product. It is well known that inequivalent unitary irreducible representations Π of \mathbb{H}^d on a Hilbert space \mathcal{H} are indexed by $\lambda \in \mathbb{R} \setminus \{0\}$ (see [10]). With them we can define the Fourier transform of $\varphi \in L^2(\mathbb{H}^d)$ as

$$\mathcal{F}_{\mathbb{H}^d} \varphi(\lambda) = \int_{\mathbb{H}^d} \varphi(x) \Pi_\lambda(x) dx \quad (1.5)$$

and those are Hilbert-Schmidt operators on \mathcal{H} .

Our definition of the bracket map for the Heisenberg group is the following.

Definition 1.1. Given $\varphi, \psi \in L^2(\mathbb{H}^d)$ and $\alpha \in (0, 1]$, we define

$$[\varphi, \psi](\alpha) = \sum_{j \in \mathbb{Z}} \langle \mathcal{F}_{\mathbb{H}^d} \varphi(\alpha + j), \mathcal{F}_{\mathbb{H}^d} \psi(\alpha + j) \rangle_{\mathcal{HS}} |\alpha + j|^d \quad (1.6)$$

where $\langle \cdot, \cdot \rangle_{\mathcal{HS}}$ denotes the inner product on the space of Hilbert-Schmidt operators on \mathcal{H} .

As we will see, the appearance of $|\alpha + j|^d$ is related to the Plancherel measure on the unitary dual of \mathbb{H}^d . Such bracket map will allow us to study PIS subspaces with respect to left translations of discrete subgroups of \mathbb{H}^d .

Left translations L are unitary representations on $L^2(\mathbb{H}^d)$ defined by

$$g \mapsto L_g \psi(x) = \psi(g^{-1}x), \quad g, x \in \mathbb{H}^d, \quad \psi \in L^2(\mathbb{H}^d) \quad (1.7)$$

and we will consider lattice subgroups Γ of \mathbb{H}^d with the following form: for $\mathbb{A}, \mathbb{B} \in GL(d, \mathbb{R})$ matrices such that $\mathbb{A}\mathbb{B}^t \in \mathbb{Z}$, we let

$$\Gamma = \mathbb{A}\mathbb{Z}^d \times \mathbb{B}\mathbb{Z}^d \times \mathbb{Z}. \quad (1.8)$$

In particular, we will distinguish the center of Γ as a group from the other variables, denoting $\Gamma_1 = \mathbb{A}\mathbb{Z}^d \times \mathbb{B}\mathbb{Z}^d$ and $\Gamma_2 = \mathbb{Z}$. This permits to write $\Gamma = \Gamma_1\Gamma_2$ in the sense that any element $\gamma \in \Gamma$ can be written as $\gamma = \gamma_1\gamma_2$, where $\gamma_i \in \Gamma_i$, $i = 1, 2$ and composition is made according to (1.4).

Properties of the bracket map (1.6) are given in Section 2, where the following result is proved.

Theorem 1.2. *Let ψ be in $L^2(\mathbb{H}^d)$, $\psi \neq 0$. Then the family $\{L_\gamma \psi\}_{\gamma \in \Gamma}$ is an orthonormal basis for $\langle \psi \rangle_{\Gamma, L}$ if and only if*

$$[\psi, L_{\gamma_1} \psi](\alpha) = \delta_{\gamma_1, 0} \quad \text{a.e. } \alpha \in (0, 1], \quad \forall \gamma_1 \in \Gamma_1. \quad (1.9)$$

Section 3 is devoted to prove a technical result of crucial importance for the results in Section 4. Namely, for $\varphi, \psi \in L^2(\mathbb{H}^d)$ we define

$$S_\psi[\varphi](\alpha) \doteq \chi_{E_\psi}(\alpha) \frac{[\varphi, \psi](\alpha)}{[\psi, \psi](\alpha)}, \quad \alpha \in (0, 1] \quad (1.10)$$

where

$$E_\psi = \{\alpha \in (0, 1] : [\psi, \psi](\alpha) > 0\}$$

and prove the following:

Theorem 1.3. *The map S_ψ is a linear isometry from*

$$\langle \psi \rangle_{\Gamma_2, L} = \overline{\text{span}\{L_{\gamma_2} \psi : \gamma_2 \in \Gamma_2\}}^{L^2(\mathbb{H}^d)}$$

onto the weighted Hilbert space $L^2((0, 1], [\psi, \psi](\alpha)d\alpha)$.

Section 4 is devoted to find conditions on ψ such that its left translations form a Riesz basis or a frame of the space $\langle \psi \rangle_{\Gamma, L}$. Such conditions in terms of the bracket map (1.6) are given by the following theorems.

Theorem 1.4. *Let ψ be in $L^2(\mathbb{H}^d)$, $\psi \neq 0$ and assume that*

$$[\psi, L_{\gamma_1} \psi](\alpha) = 0 \quad \text{a.e. } \alpha \in E_\psi, \quad \forall \gamma_1 \in \Gamma_1 \setminus \{0\}. \quad (1.11)$$

Then for $0 < A \leq B < \infty$ the following are equivalent

- a) $\{L_\gamma \psi\}_{\gamma \in \Gamma}$ is a Riesz basis for $\langle \psi \rangle_{\Gamma, L}$ with constants A and B
- b) $A \leq [\psi, \psi](\alpha) \leq B$ for a.e. $\alpha \in (0, 1]$.

Theorem 1.5. *Let ψ be in $L^2(\mathbb{H}^d)$, $\psi \neq 0$ and assume (1.11). Then for $0 < A \leq B < \infty$ the following are equivalent*

- a) $\{L_\gamma \psi\}_{\gamma \in \Gamma}$ is a frame for $\langle \psi \rangle_{\Gamma, L}$ with constants A and B
- b) $A \leq [\psi, \psi](\alpha) \leq B$ for a.e. $\alpha \in E_\psi$.

In Section 5 we will then provide a construction of frames for cyclic subspaces, based on Theorem 1.5, while Section 6 is devoted to show how the presented results can be generalized to the more abstract setting of SI/Z groups.

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2 Properties of the bracket map on the Heisenberg group

We denote by $\mathbb{H}^d \equiv \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}$ the d -dimensional Heisenberg group, with (non-commutative) composition law given by (1.4), also addressed as *polarized* Heisenberg group [10]. Note that with this law, the inverse of an element is $(p, q, t)^{-1} = (-p, -q, -t + p \cdot q)$.

The Haar measure of \mathbb{H}^d is the Lebesgue measure on \mathbb{R}^{2d+1} , and the center of \mathbb{H}^d is the subgroup $Z = \mathbb{R}$ of all elements $(0, 0, t)$, so any $x = (p, q, t)$ can be written as $x = (p, q, 0)(0, 0, t) = (0, 0, t)(p, q, 0)$.

The class of non-zero measure irreducible representations of the Heisenberg group needed to build up the Plancherel theorem are identified by $\mathbb{R}^* = \mathbb{R} \setminus \{0\}$. Then $\widehat{\mathbb{H}^d} = \mathbb{R}^*$ up to a zero measure set. For any $\lambda \in \mathbb{R}^*$ and any $x = (p, q, t) \in \mathbb{H}^d$, such representations are equivalent to Schrödinger representations on the Hilbert space $L^2(\mathbb{R}^d)$, that we denote with the map $\pi_\lambda(x)$

$$\pi_\lambda(x)f(y) = \pi_\lambda(p, q, t)f(y) = e^{2\pi i \lambda t} e^{-2\pi i \lambda q \cdot y} f(y - p) = e^{2\pi i \lambda t} \pi_\lambda(p, q, 0)f(y) \quad (2.1)$$

on $f \in L^2(\mathbb{R}^d)$, where $\pi_\lambda(p, q, 0) = M_{\lambda q} T_p$ is the shift-frequency operator on $L^2(\mathbb{R}^d)$.

For $\varphi \in L^2(\mathbb{H}^d)$ the Hilbert-Schmidt operator-valued Fourier transform $\mathcal{F}_{\mathbb{H}^d} \varphi$ is given by

$$\mathcal{F}_{\mathbb{H}^d} \varphi(\lambda) = \int_{\mathbb{H}^d} \varphi(x) \pi_\lambda(x) dx, \quad \lambda \in \mathbb{R}^*.$$

It is well known (see e.g. [10]) that the Plancherel measure on $\widehat{\mathbb{H}^d}$ is given by $|\lambda|^d d\lambda$, so that Plancherel theorem reads

$$\langle \varphi, \psi \rangle_{L^2(\mathbb{H}^d)} = \int_{\mathbb{R}} \langle \mathcal{F}_{\mathbb{H}^d} \varphi(\lambda), \mathcal{F}_{\mathbb{H}^d} \psi(\lambda) \rangle_{\mathcal{HS}} |\lambda|^d d\lambda \quad (2.2)$$

where

$$\langle \mathcal{F}_{\mathbb{H}^d} \varphi(\lambda), \mathcal{F}_{\mathbb{H}^d} \psi(\lambda) \rangle_{\mathcal{HS}} = \text{trace}_{L^2(\mathbb{R}^d)} \left(\mathcal{F}_{\mathbb{H}^d} \varphi(\lambda) (\mathcal{F}_{\mathbb{H}^d} \psi(\lambda))^\dagger \right)$$

and $(\mathcal{F}_{\mathbb{H}^d} \psi(\lambda))^\dagger$ stands for the $L^2(\mathbb{R}^d)$ adjoint of $\mathcal{F}_{\mathbb{H}^d} \psi(\lambda)$.

We will now prove some key properties of the bracket map of Definition 1.1.

Lemma 2.1. *The map (1.6) is a sesquilinear map from $L^2(\mathbb{H}^d) \times L^2(\mathbb{H}^d)$ to $L^1((0, 1])$.*

Proof. The sesquilinearity follows from the linearity of the Fourier transform and the sesquilinearity of the Hilbert-Schmidt scalar product. To show that $[\varphi, \psi] \in L^1((0, 1])$ we show that

$$\int_0^1 |[\varphi, \psi](\alpha)| d\alpha \leq \|\varphi\|_{L^2(\mathbb{H}^d)} \|\psi\|_{L^2(\mathbb{H}^d)}.$$

Indeed

$$\begin{aligned} \int_0^1 |[\varphi, \psi](\alpha)| d\alpha &= \int_0^1 \left| \sum_{j \in \mathbb{Z}} \langle \mathcal{F}_{\mathbb{H}^d} \varphi(\alpha + j), \mathcal{F}_{\mathbb{H}^d} \psi(\alpha + j) \rangle_{\mathcal{HS}} \right| |\alpha + j|^d d\alpha \\ &\leq \int_0^1 \sum_{j \in \mathbb{Z}} \|\mathcal{F}_{\mathbb{H}^d} \varphi(\alpha + j)\|_{\mathcal{HS}} \|\mathcal{F}_{\mathbb{H}^d} \psi(\alpha + j)\|_{\mathcal{HS}} |\alpha + j|^d d\alpha \\ &\leq \left(\int_0^1 \sum_{j \in \mathbb{Z}} \|\mathcal{F}_{\mathbb{H}^d} \varphi(\alpha + j)\|_{\mathcal{HS}}^2 |\alpha + j|^d d\alpha \right)^{\frac{1}{2}} \left(\int_0^1 \sum_{j \in \mathbb{Z}} \|\mathcal{F}_{\mathbb{H}^d} \psi(\alpha + j)\|_{\mathcal{HS}}^2 |\alpha + j|^d d\alpha \right)^{\frac{1}{2}} \\ &= \left(\int_{\mathbb{R}} \|\mathcal{F}_{\mathbb{H}^d} \varphi(\lambda)\|_{\mathcal{HS}}^2 |\lambda|^d d\lambda \right)^{\frac{1}{2}} \left(\int_{\mathbb{R}} \|\mathcal{F}_{\mathbb{H}^d} \psi(\lambda)\|_{\mathcal{HS}}^2 |\lambda|^d d\lambda \right)^{\frac{1}{2}} \end{aligned}$$

and by Plancherel theorem we can conclude. \square

Lemma 2.2. *For any $\varphi, \psi \in L^2(\mathbb{H}^d)$ we have*

$$|[\varphi, \psi](\alpha)|^2 \leq [\varphi, \varphi](\alpha) [\psi, \psi](\alpha) \quad \text{a.e. } \alpha \in (0, 1].$$

Proof. By Cauchy-Schwartz inequality

$$\begin{aligned} |[\varphi, \psi](\alpha)|^2 &\leq \left(\sum_{j \in \mathbb{Z}} \|\mathcal{F}_{\mathbb{H}^d} \varphi(\alpha + j)\|_{\mathcal{HS}}^2 |\alpha + j|^d \right) \left(\sum_{j \in \mathbb{Z}} \|\mathcal{F}_{\mathbb{H}^d} \psi(\alpha + j)\|_{\mathcal{HS}}^2 |\alpha + j|^d \right) \\ &= [\varphi, \varphi](\alpha) [\psi, \psi](\alpha). \end{aligned}$$

\square

The following lemma provides two properties of the bracket map that will be intensively used.

Lemma 2.3. *Let φ, ψ be in $L^2(\mathbb{H}^d)$. Then*

$$i) \quad [L_x \varphi, L_{x'} \psi](\alpha) = [\varphi, L_{x^{-1}x'} \psi](\alpha) \text{ for all } x, x' \in \mathbb{H}^d$$

$$ii) \quad [\varphi, L_{\gamma_2} \psi](\alpha) = e^{-2\pi i \gamma_2 \alpha} [\varphi, \psi](\alpha) \text{ for a.e. } \alpha \in (0, 1] \text{ and all } \gamma_2 \in \Gamma_2.$$

Proof. To prove *i)*, note first that for all $x \in \mathbb{H}^d$ (see e.g. [6])

$$\mathcal{F}_{\mathbb{H}^d} L_x \psi(\lambda) = \pi_\lambda(x) \mathcal{F}_{\mathbb{H}^d} \psi(\lambda). \quad (2.3)$$

Since the operators $\pi_\lambda(x)$, $x \in \mathbb{H}^d$ are unitary homomorphisms, then using (2.3)

$$\begin{aligned} \langle \mathcal{F}_{\mathbb{H}^d} L_x \varphi(\lambda), \mathcal{F}_{\mathbb{H}^d} L_{x'} \psi(\lambda) \rangle_{\mathcal{HS}} &= \langle \pi_\lambda(x) \mathcal{F}_{\mathbb{H}^d} \varphi(\lambda), \mathcal{F}_{\mathbb{H}^d} L_{x'} \psi(\lambda) \rangle_{\mathcal{HS}} \\ &= \langle \mathcal{F}_{\mathbb{H}^d} \varphi(\lambda), \pi_\lambda(x^{-1}) \mathcal{F}_{\mathbb{H}^d} L_{x'} \psi(\lambda) \rangle_{\mathcal{HS}} \\ &= \langle \mathcal{F}_{\mathbb{H}^d} \varphi(\lambda), \mathcal{F}_{\mathbb{H}^d} L_{x^{-1}x'} \psi(\lambda) \rangle_{\mathcal{HS}} \end{aligned}$$

so the proof follows by the definition (1.6) of the bracket.

To prove *ii*) we can use the same argument, noting that $\pi_\lambda(\gamma_2) = e^{2\pi i \gamma_2 \lambda}$. \square

Next lemma provides the key relation with the notion of dual integrability introduced in [7].

Lemma 2.4. *The abelian group Γ_2 is dual integrable in the sense of [7] on $L^2(\mathbb{H}^d)$ with respect to the bracket (1.6), i.e. for all $\gamma_2 \in \Gamma_2 = \mathbb{Z}$ we have*

$$\langle \varphi, L_{\gamma_2} \psi \rangle_{L^2(\mathbb{H}^d)} = \int_0^1 [\varphi, \psi] e^{-2\pi i \gamma_2 \alpha} d\alpha \quad \forall \varphi, \psi \in L^2(\mathbb{H}^d).$$

Proof. By Plancherel theorem and a periodization on \mathbb{R}

$$\begin{aligned} \langle \varphi, L_{\gamma_2} \psi \rangle_{L^2(\mathbb{H}^d)} &= \int_{\mathbb{R}} \langle \mathcal{F}_{\mathbb{H}^d} \varphi(\lambda), \mathcal{F}_{\mathbb{H}^d} L_{\gamma_2} \psi(\lambda) \rangle_{\mathcal{HS}} |\lambda|^d d\lambda \\ &= \int_0^1 \sum_{j \in \mathbb{Z}} \langle \mathcal{F}_{\mathbb{H}^d} \varphi(\alpha + j), \mathcal{F}_{\mathbb{H}^d} L_{\gamma_2} \psi(\alpha + j) \rangle_{\mathcal{HS}} |\alpha + j|^d d\alpha = \int_0^1 [\varphi, L_{\gamma_2} \psi] d\alpha. \end{aligned}$$

The proof then follows by Lemma 2.3, *ii*), noting that $e^{2\pi i \gamma_2 j} = 1$. \square

As a corollary of Lemma 2.4 we have the following relationship between the left regular representation of discrete subgroups (1.8) with the bracket map (1.6).

Corollary 2.5. *For all $\varphi, \psi \in L^2(\mathbb{H}^d)$ and all $\gamma = \gamma_1 \gamma_2 \in \Gamma$ we have*

$$\langle \varphi, L_\gamma \psi \rangle_{L^2(\mathbb{H}^d)} = \int_0^1 [\varphi, L_{\gamma_1} \psi] e^{-2\pi i \gamma_2 \alpha} d\alpha. \quad (2.4)$$

Proof. The proof follows by applying Lemma 2.4 to $L_{\gamma_1} \psi$ instead of ψ . \square

We can now provide an equivalent characterization of condition (1.11), that will be used in the proofs of Theorems 1.4 and 1.5.

Proposition 2.6. *Let $\psi \in L^2(\mathbb{H}^d)$. Then condition (1.11) is equivalent to*

$$\langle \psi, L_{\gamma_1 \gamma_2} \psi \rangle_{L^2(\mathbb{H}^d)} = \delta_{\gamma_1, 0} \langle \psi, L_{\gamma_2} \psi \rangle_{L^2(\mathbb{H}^d)} \quad \forall \gamma = \gamma_1 \gamma_2 \in \Gamma. \quad (2.5)$$

Proof. We first note that, by Lemma 2.2, whenever $\alpha \notin E_\psi$ then $[\psi, L_\gamma \psi] = 0$ for all γ . Hence the vanishing of the bracket in condition (1.11) can be equivalently stated for a.e. $\alpha \in (0, 1]$.

We prove the equivalence of (1.11) with (2.5) when $\gamma_1 \neq 0$, since otherwise (2.5) is always true. Let us suppose that (1.11) holds. Then (2.5) follows by (2.4) and the previous observation. On the contrary, let us suppose that (2.5) holds. Then again by (2.4) and the uniqueness of Fourier series coefficients we have that $[\psi, L_{\gamma_1} \psi] = 0$ for a.e. $\alpha \in (0, 1]$. \square

We are ready to prove the theorem concerning orthonormal systems.

Proof of Theorem 1.2. Let us first suppose that $\{L_\gamma \psi\}_{\gamma \in \Gamma}$ is an orthonormal system. Then by (2.4)

$$\delta_{\gamma,0} = \langle \psi, L_\gamma \psi \rangle_{L^2(\mathbb{H}^d)} = \int_0^1 [\psi, L_{\gamma_1} \psi](\alpha) e^{-2\pi i \alpha \gamma_2} d\alpha.$$

By Lemma 2.1, $[\psi, L_{\gamma_1} \psi]$ is in $L^1((0,1])$, so the result follows by uniqueness of Fourier series coefficients.

Conversely, let us suppose that $[\psi, L_{\gamma_1} \psi](\alpha) = \delta_{\gamma_1,0}$ for a.e. $\alpha \in (0,1]$ and all $\gamma_1 \in \Gamma_1$. Then if we call $\gamma'' = \gamma^{-1}\gamma'$, and decompose it as $\gamma'' = \gamma_1''\gamma_2''$ with $\gamma_i'' \in \Gamma_i, i = 1, 2$

$$\begin{aligned} \langle L_\gamma \psi, L_{\gamma'} \psi \rangle_{L^2(\mathbb{H}^d)} &= \langle \psi, L_{\gamma^{-1}\gamma'} \psi \rangle_{L^2(\mathbb{H}^d)} = \int_0^1 [\psi, L_{\gamma_1''} \psi](\alpha) e^{-2\pi i \alpha \gamma_2''} d\alpha \\ &= \delta_{\gamma_1'',0} \int_0^1 e^{-2\pi i \alpha \gamma_2''} d\alpha = \delta_{\gamma_1'',0} \delta_{\gamma_2'',0} \end{aligned}$$

so the result follows since $\delta_{\gamma_1'',0} \delta_{\gamma_2'',0} = \delta_{\gamma,\gamma'}$. \square

3 An isometric isomorphism

Let ψ be in $L^2(\mathbb{H}^d)$, and call $E_\psi := \{\alpha \in \mathbb{T} : [\psi, \psi](\alpha) > 0\}$. Recall that the map S_ψ at $\varphi \in L^2(\mathbb{H}^d)$ was defined in (1.10) as

$$S_\psi[\varphi](\alpha) := \chi_{E_\psi}(\alpha) \frac{[\varphi, \psi](\alpha)}{[\psi, \psi](\alpha)} \quad \text{a.e. } \alpha \in (0,1].$$

The following lemma shows how S_ψ acts on Γ shifts of a function.

Lemma 3.1. *Let ψ be in $L^2(\mathbb{H}^d)$. Then*

i) for all $\varphi \in L^2(\mathbb{H}^d)$ and all $\Gamma \ni \gamma = \gamma_1 \gamma_2$ with $\gamma_i \in \Gamma_i, i = 1, 2$

$$S_\psi[L_\gamma \varphi](\alpha) = e^{2\pi i \gamma_2 \alpha} S_\psi[L_{\gamma_1} \varphi](\alpha) \quad \text{a.e. } \alpha \in (0,1]$$

ii) if $\varphi = \sum_{\gamma_2 \in \Gamma_2} a_{\gamma_2} L_{\gamma_2} \psi$ be a finite sum with complex coefficients a_{γ_2} , then

$$S_\psi[\varphi](\alpha) = S_\psi \left[\sum_{\gamma_2 \in \Gamma_2} a_{\gamma_2} L_{\gamma_2} \psi \right](\alpha) = \left(\sum_{\gamma_2 \in \Gamma_2} a_{\gamma_2} e^{2\pi i \gamma_2 \alpha} \right) \chi_{E_\psi}(\alpha) \quad \text{a.e. } \alpha \in (0,1].$$

Proof. The proof of i) follows from Lemma 2.3, ii). To prove ii), we see that i) extends to finite sums by the linearity of the bracket map with respect to its first argument. Moreover, for γ_2 shifts we end up with $S_\psi[\varphi](\alpha)$, equal to $\chi_{E_\psi}(\alpha)$. \square

Proof of Theorem 1.3. The linearity of the map follows from its definition and the linearity of the bracket function $[\cdot, \cdot]$ on its first argument. Let $\varphi \in L^2(\mathbb{H}^d)$. Using the definition of S_ψ and Lemma 2.2 we have

$$\begin{aligned} \int_0^1 |S_\psi[\varphi](\alpha)|^2 [\psi, \psi](\alpha) d\alpha &= \int_{E_\psi} \frac{|[\varphi, \psi](\alpha)|^2}{[\psi, \psi](\alpha)} d\alpha \leq \int_{E_\psi} \frac{[\varphi, \varphi](\alpha)[\psi, \psi](\alpha)}{[\psi, \psi](\alpha)} d\alpha \quad (3.1) \\ &= \int_0^1 [\varphi, \varphi](\alpha) d\alpha = \|\varphi\|^2 \end{aligned}$$

where the last identity is due to Lemma 2.4. We have thus proved that the image of the map S_ψ is in the Hilbert space $L^2(\mathbb{T}, [\psi, \psi](\alpha) d\alpha)$.

To prove that the restriction of S_ψ to the space $\langle \psi \rangle_{\Gamma_2, L}$ is an isometry (hence one-to-one) onto $L^2((0, 1], [\psi, \psi](\alpha) d\alpha)$, we apply Lemma 3.1 as follows. If $\varphi \in \text{span}\{L_{\gamma_2}\psi : \gamma_2 \in \Gamma_2\}$, then we can write $\varphi = \sum_{\gamma_2 \in \Lambda} a_{\gamma_2} L_{\gamma_2}\psi$ for some finite sequence $\{a_{\gamma_2}\}_{\gamma_2 \in \Lambda \subset \Gamma_2}$. Now by the sesquilinearity of S_ψ and using Lemma 2.3, *ii*)

$$\begin{aligned} \int_0^1 \left| S \left[\sum_{\gamma_2 \in \Lambda} a_{\gamma_2} L_{\gamma_2}\psi \right] (\alpha) \right|^2 [\psi, \psi](\alpha) d\alpha &= \int_0^1 \left| \sum_{\gamma_2} a_{\gamma_2} e^{2\pi i k \alpha} \right|^2 [\psi, \psi](\alpha) d\alpha \\ &= \int_0^1 \left[\sum_{\gamma_2} a_{\gamma_2} L_{\gamma_2}\psi, \sum_{\gamma_2} a_{\gamma_2} L_{\gamma_2}\psi \right] (\alpha) d\alpha = \int_0^1 [\varphi, \varphi](\alpha) d\alpha = \|\varphi\|^2. \end{aligned}$$

For general $\varphi \in \langle \psi \rangle_{\Gamma_2, L}$ we have that for any $\epsilon > 0$ we can choose a finite $\Lambda_\epsilon \subset \Gamma_2$ and $\varphi_\epsilon = \sum_{\gamma_2 \in \Lambda_\epsilon} a_{\gamma_2} L_{\gamma_2}\psi$ such that $\|\varphi - \varphi_\epsilon\| \leq \epsilon$. Then

$$\|\varphi\| \leq \|\varphi - \varphi_\epsilon\| + \|\varphi_\epsilon\| \leq \epsilon + \|\varphi_\epsilon\|. \quad (3.2)$$

Moreover

$$\begin{aligned} \|\varphi_\epsilon\| &= \left(\int_0^1 |S_\psi[\varphi_\epsilon](\alpha)|^2 [\psi, \psi](\alpha) d\alpha \right)^{\frac{1}{2}} \\ &\leq \left(\int_0^1 |S_\psi[\varphi_\epsilon - \varphi](\alpha)|^2 [\psi, \psi](\alpha) d\alpha \right)^{\frac{1}{2}} + \left(\int_0^1 |S_\psi[\varphi](\alpha)|^2 [\psi, \psi](\alpha) d\alpha \right)^{\frac{1}{2}} \\ &\leq \|\varphi - \varphi_\epsilon\| + \left(\int_0^1 |S_\psi[\varphi](\alpha)|^2 [\psi, \psi](\alpha) d\alpha \right)^{\frac{1}{2}} \end{aligned}$$

where the last inequality is due to (3.1). By (3.2) we can then deduce that

$$\left(\int_0^1 |S_\psi[\varphi](\alpha)|^2 [\psi, \psi](\alpha) d\alpha \right)^{\frac{1}{2}} \geq \|\varphi_\epsilon\| - \|\varphi - \varphi_\epsilon\| \geq \|\varphi_\epsilon\| - \epsilon \geq \|\varphi\| - 2\epsilon.$$

This combined with (3.1) provides then the isometry for the space $\langle \psi \rangle_{\Gamma_2, L}$.

To show that S_ψ is onto we argue by contradiction. Suppose that $S_\psi(\langle\psi\rangle_{\Gamma_2,L}) \subsetneq L^2((0,1], [\psi, \psi](\alpha)d\alpha)$ holds. Then we can choose a nontrivial $f \in L^2((0,1], [\psi, \psi](\alpha)d\alpha)$ such that $f \perp S_\psi(\langle\psi\rangle_{\Gamma_2,L})$, i.e.

$$0 = \int_0^1 f(\alpha) \overline{S_\psi[L_{\gamma_2}\psi](\alpha)} [\psi, \psi](\alpha) d\alpha \quad \forall \gamma_2 \in \Gamma_2.$$

But since $S_\psi[L_{\gamma_2}\psi](\alpha) = e^{2\pi i\alpha\gamma_2} S_\psi[\psi](\alpha) = e^{2\pi i\alpha\gamma_2} \chi_{E_\psi}$ we are saying that

$$0 = \int_0^1 e^{-2\pi i\alpha\gamma_2} f(\alpha) [\psi, \psi](\alpha) d\alpha \quad \forall \gamma_2 \in \Gamma_2$$

so, by uniqueness of Fourier series coefficients and the nonnegativity of $[\psi, \psi](\alpha)$, this implies that $f = 0$ in $L^2((0,1], [\psi, \psi](\alpha)d\alpha)$, hence the contradiction. \square

4 Riesz bases and frames

In this section we prove Theorems 1.4 and 1.5. Recall that $\{L_\gamma\psi\}_{\gamma \in \Gamma}$ is a Riesz basis for $\langle\psi\rangle_{\Gamma,L}$ with constants A and B if it is a basis and if for all sequences $\{a_\gamma\}_\Gamma \in \ell_2(\Gamma)$ it satisfies

$$A \sum_{\gamma \in \Gamma} |a_\gamma|^2 \leq \left\| \sum_{\gamma \in \Gamma} a_\gamma L_\gamma \psi \right\|_{L^2(\mathbb{H}^d)}^2 \leq B \sum_{\gamma \in \Gamma} |a_\gamma|^2, \quad (4.1)$$

while $\{L_\gamma\psi\}_{\gamma \in \Gamma}$ is a frame for $\langle\psi\rangle_{\Gamma,L}$ with constants A and B if for every $\varphi \in \langle\psi\rangle_{\Gamma,L}$ it holds

$$A \|\varphi\|_{L^2(\mathbb{H}^d)}^2 \leq \sum_{\gamma \in \Gamma} |\langle \varphi, L_\gamma \psi \rangle_{L^2(\mathbb{H}^d)}|^2 \leq B \|\varphi\|_{L^2(\mathbb{H}^d)}^2. \quad (4.2)$$

Both theorems require condition (1.11) or its equivalent forms described by Proposition 2.6. One relevant consequence of such condition is the following.

Lemma 4.1. *Let $\psi \in L^2(\mathbb{H}^d)$ and assume condition (1.11). Then,*

i) *If $\gamma_1, \eta_1 \in \Gamma_1$ and $\gamma_1 \neq \eta_1$, $\langle L_{\gamma_1}\psi \rangle_{\Gamma_2,L} \perp \langle L_{\eta_1}\psi \rangle_{\Gamma_2,L}$ and*

$$\langle\psi\rangle_{\Gamma,L} = \bigoplus_{\gamma_1 \in \Gamma_1} \langle L_{\gamma_1}\psi \rangle_{\Gamma_2,L}. \quad (4.3)$$

as an orthogonal direct sum.

ii) *For each $\varphi \in \langle\psi\rangle_{\Gamma,L}$ there exist unique $\varphi_{\gamma_1} \in \langle L_{\gamma_1}\psi \rangle_{\Gamma_2,L}$, $\gamma_1 \in \Gamma_1$, such that $\varphi = \sum_{\gamma_1 \in \Gamma_1} \varphi_{\gamma_1}$ (convergence in $L^2(\mathbb{H}^d)$) and*

$$\|\varphi\|_{L^2(\mathbb{H}^d)}^2 = \sum_{\gamma_1 \in \Gamma_1} \|\varphi_{\gamma_1}\|_{L^2(\mathbb{H}^d)}^2. \quad (4.4)$$

Proof. Follows from Proposition 2.6 by easy density arguments. \square

Proposition 4.2. *Let ψ be in $L^2(\mathbb{H}^d)$ and assume condition (1.11). Then*

i) *$\{L_\gamma\psi\}_{\gamma \in \Gamma}$ is a Riesz family with constants A and B if and only if $\{L_{\gamma_2}\psi\}_{\gamma_2 \in \Gamma_2}$ is a Riesz family with constants A and B .*

ii) *$\{L_\gamma\psi\}_{\gamma \in \Gamma}$ is a frame for $\langle\psi\rangle_{\Gamma,L}$ with constants A and B if and only if $\{L_{\gamma_2}\psi\}_{\gamma_2 \in \Gamma_2}$ is a frame for $\langle\psi\rangle_{\Gamma_2,L}$ with constants A and B .*

Proof. i) Suppose that $\{L_{\gamma_2}\psi\}_{\gamma_2 \in \Gamma_2}$ is a Riesz family with constants A and B . Let $\{a_\gamma\}_{\gamma \in \Gamma} \in \ell_2(\Gamma)$. Write

$$\varphi \doteq \sum_{\gamma \in \Gamma} a_\gamma L_\gamma \psi = \sum_{\gamma_1} \sum_{\gamma_2} a_{\gamma_1 \gamma_2} L_{\gamma_2} L_{\gamma_1} \psi \doteq \sum_{\gamma_1} \varphi_{\gamma_1},$$

with $\varphi_{\gamma_1} \doteq \sum_{\gamma_2} a_{\gamma_1 \gamma_2} L_{\gamma_2} L_{\gamma_1} \psi \in \langle L_{\gamma_1} \psi \rangle_{\Gamma_2, L}$ since $\{L_{\gamma_2} \psi\}_{\gamma_2 \in \Gamma_2}$ is a Riesz family and $\{a_{\gamma_1 \gamma_2}\}_{\gamma_2 \in \Gamma_2} \in \ell_2(\Gamma_2)$ for each $\gamma_1 \in \Gamma_1$. By (4.4)

$$\begin{aligned} \left\| \sum_{\gamma \in \Gamma} a_\gamma L_\gamma \psi \right\|_{L^2(\mathbb{H}^d)}^2 &= \sum_{\gamma_1} \left\| \sum_{\gamma_2} a_{\gamma_1 \gamma_2} L_{\gamma_2} L_{\gamma_1} \psi \right\|_{L^2(\mathbb{H}^d)}^2 \\ &= \sum_{\gamma_1} \left\| \sum_{\gamma_2} a_{\gamma_1 \gamma_2} L_{\gamma_2} \psi \right\|_{L^2(\mathbb{H}^d)}^2 \end{aligned}$$

since L_{γ_1} is an isometry on $L^2(\mathbb{H}^d)$. Our hypothesis implies

$$\left\| \sum_{\gamma \in \Gamma} a_\gamma L_\gamma \psi \right\|_{L^2(\mathbb{H}^d)}^2 \approx \sum_{\gamma_1} \sum_{\gamma_2} |a_{\gamma_1 \gamma_2}|^2 = \sum_{\gamma \in \Gamma} |a_\gamma|^2,$$

where the equivalence is up to the Riesz family constants A and B . This shows (4.1) for the family $\{L_\gamma \psi\}_{\gamma \in \Gamma}$. The converse is straight forward.

ii) Suppose that $\{L_{\gamma_2} \psi\}_{\gamma_2 \in \Gamma_2}$ is a frame for $\langle \psi \rangle_{\Gamma_2, L}$ with constants A and B . By Lemma 4.1 ii), choose $\varphi_{\gamma_1} \in \langle L_{\gamma_1} \psi \rangle_{\Gamma_2, L}$, $\gamma_1 \in \Gamma_1$, such that $\varphi = \sum_{\gamma_1} \varphi_{\gamma_1}$ (convergence in $L^2(\mathbb{H}^d)$) and (4.4) holds. Then, by the orthogonality relations in Lemma 4.1 i),

$$\begin{aligned} \sum_{\gamma \in \Gamma} |\langle \varphi, L_\gamma \psi \rangle_{L^2(\mathbb{H}^d)}|^2 &= \sum_{\eta_1} \sum_{\eta_2} \left| \langle \sum_{\gamma_1} \varphi_{\gamma_1}, L_{\eta_1} L_{\eta_2} \psi \rangle_{L^2(\mathbb{H}^d)} \right|^2 \\ &= \sum_{\gamma_1} \sum_{\eta_2} |\langle \varphi_{\gamma_1}, L_{\gamma_1} L_{\eta_2} \psi \rangle_{L^2(\mathbb{H}^d)}|^2 \\ &= \sum_{\gamma_1} \sum_{\eta_2} |\langle L_{\gamma_1}^{-1} \varphi_{\gamma_1}, L_{\eta_2} \psi \rangle_{L^2(\mathbb{H}^d)}|^2. \end{aligned}$$

Since $L_{\gamma_1}^{-1} \varphi_{\gamma_1} \in \langle \psi \rangle_{\Gamma_2, L}$ our hypothesis implies

$$\sum_{\gamma \in \Gamma} |\langle \varphi, L_\gamma \psi \rangle_{L^2(\mathbb{H}^d)}|^2 \approx \sum_{\gamma_1} \|L_{\gamma_1}^{-1} \varphi_{\gamma_1}\|_{L^2(\mathbb{H}^d)}^2 = \sum_{\gamma_1} \|\varphi_{\gamma_1}\|_{L^2(\mathbb{H}^d)}^2 = \|\varphi\|_{L^2(\mathbb{H}^d)}^2,$$

by (4.4) where the equivalence is up to the frame constants A and B . This proves that $\{L_\gamma \psi\}_{\gamma \in \Gamma}$ is a frame for $\langle \psi \rangle_{\Gamma, L}$ with constants A and B .

For the converse, let $\varphi \in \langle \psi \rangle_{\Gamma_2, L}$. By the orthogonality relation in Lemma 4.1 i), $\langle \varphi, L_{\gamma_1} L_{\gamma_2} \psi \rangle = 0$ for all $\gamma_1 \in \Gamma_1$, $\gamma_1 \neq 0$. Then

$$\sum_{\gamma_2} |\langle \varphi, L_{\gamma_2} \psi \rangle_{L^2(\mathbb{H}^d)}|^2 = \sum_{\gamma_1} \sum_{\gamma_2} |\langle \varphi, L_{\gamma_1} L_{\gamma_2} \psi \rangle_{L^2(\mathbb{H}^d)}|^2 \approx \|\varphi\|_{L^2(\mathbb{H}^d)}^2,$$

since $\varphi \in \langle \psi \rangle_{\Gamma_2, L} \subset \langle \psi \rangle_{\Gamma, L}$, where the equivalence is up to the frame constants A and B . \square

We are now ready to prove the two main theorems.

Proof of Theorem 1.4. Let us first assume *b*). Observe that by Lemma 2.3, *i*), *b*) is equivalent to

$$A \leq [L_{\gamma_1} \psi, L_{\gamma_1} \psi](\alpha) \leq B \quad \text{a.e. } \alpha \in (0, 1], \forall \gamma_1 \in \Gamma_1.$$

This implies that $E_{\psi_{\gamma_1}} = E_\psi = (0, 1]$ a.e. Hence, by Lemma (3.1) *i*)

$$S_{L_{\gamma_1} \psi}[L_{\gamma_2} L_{\gamma_1} \psi](\alpha) = e^{2\pi i \alpha \gamma_2} \quad \text{a.e. } \alpha \in (0, 1].$$

Moreover, we have that $L^2((0, 1], [L_{\gamma_1} \psi, L_{\gamma_1} \psi](\alpha) d\alpha) \approx L^2((0, 1], d\alpha)$ for all $\gamma_1 \in \Gamma_1$, where \approx stands for equivalence up to constants A and B . Since, by Theorem 1.3, $S_{L_{\gamma_1} \psi}$ is a linear isometry from $\langle L_{\gamma_1} \psi \rangle_{\Gamma_2, L}$ onto $L^2((0, 1], [L_{\gamma_1} \psi, L_{\gamma_1} \psi](\alpha) d\alpha)$, and since $\{e^{2\pi i \alpha \gamma_2}\}_{\gamma_2 \in \Gamma_2}$ is an orthonormal basis of $L^2((0, 1], d\alpha)$, we conclude that $\{L_{\gamma_2} L_{\gamma_1} \psi\}_{\gamma_2 \in \Gamma_2}$ is a basis of $\langle L_{\gamma_1} \psi \rangle_{\Gamma_2, L}$ for all $\gamma_1 \in \Gamma_1$. Thus by (4.3) we have that $\{L_\gamma \psi\}_{\gamma \in \Gamma}$ is a basis of $\langle \psi \rangle_{\Gamma, L}$.

Now, using Proposition 4.2 *i*), to conclude the proof that *b*) implies *a*) and to prove that *a*) implies *b*) we need only to prove that the following conditions are equivalent:

a') $\{L_{\gamma_2} \psi\}_{\gamma_2 \in \Gamma_2}$ is a Riesz family for $\langle \psi \rangle_{\Gamma_2, L}$ with constants A and B

b) $A \leq [\psi, \psi](\alpha) \leq B$ for a.e. $\alpha \in (0, 1]$.

If *b*) holds, and $\{c_{\gamma_2}\}_{\gamma_2 \in \Gamma_2} \in \ell_2(\Gamma_2)$, by Theorem 1.3 and Lemma 3.1,

$$\begin{aligned} \left\| \sum_{\gamma_2} c_{\gamma_2} L_{\gamma_2} \psi \right\|_{L^2(\mathbb{H}^d)}^2 &= \left\| S_\psi \left[\sum_{\gamma_2} c_{\gamma_2} L_{\gamma_2} \psi \right] \right\|_{L^2((0, 1], [\psi, \psi](\alpha) d\alpha)}^2 \\ &\approx \left\| S_\psi \left[\sum_{\gamma_2} c_{\gamma_2} L_{\gamma_2} \psi \right] \right\|_{L^2((0, 1], d\alpha)}^2 \\ &= \int_0^1 \left| \sum_{\gamma_2} c_{\gamma_2} e^{2\pi i \gamma_2 \alpha} \right|^2 d\alpha = \sum_{\gamma_2} |c_{\gamma_2}|^2, \end{aligned}$$

where equivalence is up to constants A and B . Thus we have obtained *a')*.

Conversely, let us suppose that *a')* holds and proceed by contradiction. Let us assume that $[\psi, \psi](\alpha) < A$ on a set $E \subset (0, 1]$ of positive measure. Since $\chi_E \in L^2((0, 1], d\alpha)$, there exists a sequence $\{c_{\gamma_2}\}_{\gamma_2 \in \Gamma_2} \in \ell_2(\Gamma_2)$ such that

$$\chi_E(\alpha) = \sum_{\gamma_2} c_{\gamma_2} e^{2\pi i \gamma_2 \alpha}, \quad \text{and } |E| = \int_0^1 |\chi_E(\alpha)|^2 d\alpha = \sum_{\gamma_2} |c_{\gamma_2}|^2.$$

But, again using Theorem 1.3 and Lemma 3.1, *ii*)

$$\begin{aligned} \left\| \sum_{\gamma_2} c_{\gamma_2} L_{\gamma_2} \psi \right\|_{L^2(\mathbb{H}^d)}^2 &= \int_0^1 \left| S_\psi \left[\sum_{\gamma_2} c_{\gamma_2} L_{\gamma_2} \psi \right](\alpha) \right|^2 [\psi, \psi](\alpha) d\alpha \\ &= \int_0^1 |\chi_E(\alpha)|^2 [\psi, \psi](\alpha) d\alpha < A|E| = A \sum_{\gamma_2} |c_{\gamma_2}|^2 \end{aligned}$$

which contradicts *a')*. This proves that $[\psi, \psi](\alpha) \geq A$ for a.e. $\alpha \in (0, 1]$. The proof that $[\psi, \psi](\alpha) \leq B$ for a.e. $\alpha \in (0, 1]$ is similar. \square

Proof of Theorem 1.5. As in the previous proof, by Proposition 4.2 it suffices to prove that $b)$ is equivalent to

$a')$ $\{L_{\gamma_2}\psi\}_{\gamma_2 \in \Gamma_2}$ is a frame for $\langle \psi \rangle_{\Gamma_2, L}$ with constants A and B .

Let us first assume $b)$, and choose a $\varphi \in \langle \psi \rangle_{\Gamma_2, L}$. Then by Theorem 1.3 and Lemma 3.1, $i)$

$$\begin{aligned} \sum_{\gamma_2 \in \Gamma_2} |\langle \varphi, L_{\gamma_2}\psi \rangle_{L^2(\mathbb{H}^d)}|^2 &= \sum_{\gamma_2 \in \Gamma_2} |\langle S_\psi[\varphi], S_\psi[L_{\gamma_2}\psi] \rangle_{L^2((0,1], [\psi, \psi](\alpha)d\alpha)}|^2 \\ &= \sum_{\gamma_2 \in \Gamma_2} \left| \int_0^1 S_\psi[\varphi](\alpha) e^{-2\pi i \alpha \gamma_2} [\psi, \psi](\alpha) d\alpha \right|^2 \\ &= \int_0^1 |S_\psi[\varphi](\alpha)|^2 |[\psi, \psi](\alpha)|^2 d\alpha \end{aligned} \quad (4.5)$$

and the last transition is due to Parseval's identity. So if $A \leq [\psi, \psi](\alpha) \leq B$ for a.e. $\alpha \in E_\psi$, we get

$$\sum_{\gamma_2 \in \Gamma_2} |\langle \varphi, L_{\gamma_2}\psi \rangle_{L^2(\mathbb{H}^d)}|^2 \approx \int_0^1 |S_\psi[\varphi](\alpha)|^2 [\psi, \psi](\alpha) d\alpha = \|\varphi\|_{L^2(\mathbb{H}^d)}^2$$

where the last identity is due to Theorem 1.3 and equivalence is up to constants A and B .

Now let us assume $a')$ and suppose by contradiction that $[\psi, \psi](\alpha) < A$ on a set $E \subset E_\psi$ of positive measure. Observe that $\chi_E \in L^2((0,1], [\psi, \psi](\alpha)d\alpha)$ because $[\psi, \psi] \in L^1((0,1], d\alpha)$. Since S_ψ is onto, we can then choose $\varphi_E \in \langle \psi \rangle_{\Gamma_2, L}$ such that $S_\psi[\varphi_E] = \chi_E$. As for (4.5)

$$\begin{aligned} \sum_{\gamma_2 \in \Gamma_2} |\langle \varphi_E, L_{\gamma_2}\psi \rangle_{L^2(\mathbb{H}^d)}|^2 &= \int_0^1 |S_\psi[\varphi_E](\alpha)|^2 |[\psi, \psi](\alpha)|^2 d\alpha \\ &= \int_0^1 |\chi_E(\alpha)|^2 |[\psi, \psi](\alpha)|^2 d\alpha < A \int_E [\psi, \psi](\alpha) d\alpha \end{aligned}$$

but since S_ψ is an isometry

$$\int_E [\psi, \psi](\alpha) d\alpha = \int_0^1 |S_\psi[\varphi_E](\alpha)|^2 [\psi, \psi](\alpha) d\alpha = \|\varphi_E\|_{L^2(\mathbb{H}^d)}^2,$$

hence contradicting $a')$. This proves that $[\psi, \psi](\alpha) \geq A$ for a.e. $\alpha \in E_\psi$. The proof that $[\psi, \psi](\alpha) \leq B$ for a.e. $\alpha \in E_\psi$ is similar. \square

5 Applications

In this section we show as an application of the previous theory how frames for cyclic subspaces can be associated to orthonormal Gabor systems $\in L^2(\mathbb{R}^d)$.

Let u be a unit norm element in $L^2(\mathbb{R}^d)$, and call u_λ its normalized scalings

$$u_\lambda(x) \doteq |\lambda|^{\frac{d}{2}} u(\lambda x), \quad \lambda \in \mathbb{R}^*.$$

Let \mathbb{P}_λ be the associated projector on $L^2(\mathbb{R}^d)$

$$\mathbb{P}_\lambda = u_\lambda \otimes u_\lambda, \quad \mathbb{P}_\lambda : L^2(\mathbb{R}^d) \ni f \mapsto \langle f, u_\lambda \rangle_{L^2(\mathbb{R}^d)} u_\lambda$$

and define, for all $\epsilon \in (0, 1)$

$$H_\epsilon(\lambda) = \begin{cases} \mathbb{P}_\lambda & \epsilon < \lambda \leq 1 \\ 0 & \text{otherwise} \end{cases}.$$

For all $\epsilon \in (0, 1)$, H_ϵ belongs to $L^2(\mathbb{R}^*, \mathcal{HS}(L^2(\mathbb{R}^d)), |\lambda|^d d\lambda)$, since

$$\|H_\epsilon(\lambda)\|_{\mathcal{HS}(L^2(\mathbb{R}^d))} = \begin{cases} 1 & \epsilon < \lambda \leq 1 \\ 0 & \text{otherwise} \end{cases}$$

so

$$\|H_\epsilon(\lambda)\|_{L^2(\mathbb{R}^*, \mathcal{HS}(L^2(\mathbb{R}^d)), |\lambda|^d d\lambda)}^2 = \int_{\mathbb{R}} \|H_\epsilon(\lambda)\|_{\mathcal{HS}(L^2(\mathbb{R}^d))}^2 |\lambda|^d d\lambda = \int_{\epsilon}^1 \lambda^d d\lambda = \frac{1 - \epsilon^{d+1}}{d+1}.$$

We can then define $\psi_\epsilon \in L^2(\mathbb{H}^d)$ as $\psi_\epsilon = \mathcal{F}_{\mathbb{H}^d}^{-1} H_\epsilon$, and prove the following theorem.

Theorem 5.1. *Let $\{\pi_\alpha(\gamma_1)u_\alpha\}_{\gamma_1 \in \Gamma_1}$ be an orthonormal system in $L^2(\mathbb{R}^d)$ for a.e. $0 < \epsilon < \alpha \leq 1$. Then $\{L_\gamma \psi_\epsilon\}$ is a frame for $\langle \psi_\epsilon \rangle_{\Gamma, L}$ for all $\epsilon \in (0, 1)$.*

In order to prove Theorem 5.1 we start with two preliminary lemmata.

Lemma 5.2. *For all $\epsilon \in (0, 1)$ and all $\alpha \in E_{\psi_\epsilon}$*

$$\epsilon^d \leq [\psi_\epsilon, \psi_\epsilon](\alpha) \leq 1.$$

Proof. Since ψ_ϵ is bandlimited

$$\begin{aligned} [\psi_\epsilon, \psi_\epsilon](\alpha) &= \sum_{j \in \mathbb{Z}} \|\mathcal{F}_{\mathbb{H}^d} \psi_\epsilon(\alpha + j)\|_{\mathcal{HS}(L^2(\mathbb{R}^d))}^2 |\alpha + j|^d \\ &= \sum_{j \in \mathbb{Z}} \|H_\epsilon(\alpha + j)\|_{\mathcal{HS}(L^2(\mathbb{R}^d))}^2 |\alpha + j|^d \\ &= \|H_\epsilon(\alpha)\|_{\mathcal{HS}(L^2(\mathbb{R}^d))} |\alpha|^d = \begin{cases} \alpha^d & \epsilon < \alpha \leq 1 \\ 0 & \text{otherwise} \end{cases}. \end{aligned}$$

□

Lemma 5.3. *For all $\epsilon \in (0, 1)$, all $\gamma_1 \in \Gamma_1$ and all $\alpha \in E_{\psi_\epsilon}$*

$$[L_{\gamma_1} \psi_\epsilon, \psi_\epsilon](\alpha) = \langle \pi_\alpha(\gamma_1)u_\alpha, u_\alpha \rangle_{L^2(\mathbb{R}^d)}.$$

Proof. Again using the bandlimitness of ψ_ϵ to reduce the series to the only nonzero term

$$\begin{aligned} [L_{\gamma_1} \psi_\epsilon, \psi_\epsilon](\alpha) &= \sum_{j \in \mathbb{Z}} \langle \mathcal{F}_{\mathbb{H}^d} L_{\gamma_1} \psi_\epsilon(\alpha + j), \mathcal{F}_{\mathbb{H}^d} \psi_\epsilon(\alpha + j) \rangle_{\mathcal{HS}(L^2(\mathbb{R}^d))} |\alpha + j|^d \\ &= \sum_{j \in \mathbb{Z}} \langle \pi_{\alpha+j}(\gamma_1) H_\epsilon(\alpha + j), H_\epsilon(\alpha + j) \rangle_{\mathcal{HS}(L^2(\mathbb{R}^d))} |\alpha + j|^d \\ &= \langle \pi_\alpha(\gamma_1) H_\epsilon(\alpha), H_\epsilon(\alpha) \rangle_{\mathcal{HS}(L^2(\mathbb{R}^d))} |\alpha|^d. \end{aligned}$$

Now, since $\pi_\alpha(\gamma_1)\mathbb{P}_\alpha = (\pi_\alpha(\gamma_1)u_\alpha) \otimes u_\alpha$, when $\epsilon < \alpha \leq 1$ we have

$$\begin{aligned} \langle \pi_\alpha(\gamma_1)H_\epsilon(\alpha), H_\epsilon(\alpha) \rangle_{\mathcal{HS}(L^2(\mathbb{R}^d))} &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \overline{u_\alpha(x)} (\pi_\alpha(\gamma_1)u_\alpha(y)) \overline{u_\alpha(y)} u_\alpha(x) dx dy \\ &= \langle \pi_\alpha(\gamma_1)u_\alpha, u_\alpha \rangle_{L^2(\mathbb{R}^d)}. \end{aligned}$$

□

Proof of Theorem 5.1. In order to get a frame, we need only to verify condition b) of Theorem 1.5, which is granted by Lemma 5.2, and condition (1.11), that holds by Lemma 5.3 whenever $\{\pi_\alpha(\gamma_1)u_\alpha\}_{\gamma_1 \in \Gamma_1}$ is an orthonormal system. □

Example 5.4. One example of unitary $u \in L^2(\mathbb{R}^d)$ and Γ_1 for which $\{\pi_\alpha(\gamma_1)u_\alpha\}_{\gamma_1 \in \Gamma_1}$ is an orthonormal system of $L^2(\mathbb{R}^d)$ for $0 < \epsilon < \alpha \leq 1$ is the following. Choose $u = \chi_{(0,1)} \times \cdots \times \chi_{(0,1)} \in L^2(\mathbb{R}^d)$ and $\Gamma_1 = a\mathbb{Z}^d \times b\mathbb{Z}^d$ with $a, b \in \mathbb{Z}$. The desired orthonormality follows from

$$\pi_\alpha(\gamma_1)u_\alpha(y) = |\alpha|^{d/2} e^{-2\pi i \alpha b \langle m, y \rangle} \chi_{(0, \frac{1}{\alpha}) \times \cdots \times (0, \frac{1}{\alpha})}(y - an)$$

for $\gamma_1 = (an, bm, 0) \in \Gamma_1$.

6 SI/Z groups

In this section we consider a class of nilpotent groups distinct from the stratified groups: following [2], we say that a simply connected nilpotent Lie group \mathcal{G} is an *SI/Z group* if almost all of its irreducible representations Π are square-integrable modulo the center of the group. An irreducible representation Π of \mathcal{G} on the Hilbert space \mathcal{H} is *square-integrable modulo the center* if there exist $u, v \in \mathcal{H}$ such that

$$\int_{\mathcal{G}/Z} |\langle \Pi(n) u, v \rangle_{\mathcal{H}}|^2 d\dot{n} < \infty.$$

We remark that the class of SI/Z groups is broad and in particular contains the Heisenberg group as well as groups of an arbitrarily high degree of non-commutativity.

The effect of the SI/Z condition is that for this class of groups, the operator-valued Plancherel transform retains certain key features of the Euclidean case, and in particular, makes it possible to extend to SI/Z groups the notion of bracket map. The following construction, leading to the Plancherel theorem, can be found in [3, 2].

For a SI/Z group \mathcal{G} , fix a basis $\{X_1, \dots, X_n\}$ for its Lie algebra \mathfrak{g} , ordered in such a way that $\mathfrak{g}_j = \text{span}\{X_1, \dots, X_j\}$ is an ideal in \mathfrak{g} (central series) and \mathfrak{g}_r is the center \mathfrak{z} of \mathfrak{g} for some $r = n - 2d$. Then the center Z of \mathcal{G} can be identified with \mathbb{R}^r by choosing exponential coordinates

$$Z = e^{z_1 X_1} e^{z_2 X_2} \dots e^{z_r X_r}$$

where $(z_1, z_2, \dots, z_r) \in \mathbb{R}^r$, while the subset of \mathcal{G}

$$\mathcal{X} = e^{\mathbb{R}X_{r+1}} e^{\mathbb{R}X_{r+2}} \dots e^{\mathbb{R}X_n}$$

is identified by \mathbb{R}^{2d} . For any $g \in \mathcal{G}$ there is a unique $x \in \mathcal{X}$ and a unique $z \in Z$ such that $g = xz = zx$, therefore the group \mathcal{G} is identified by $\mathbb{R}^{2d} \times \mathbb{R}^r$.

Let \mathfrak{g}^* denote the linear dual of \mathfrak{g} , and let \mathfrak{z}^* denote the subspace

$$\mathfrak{z}^* = \{\lambda \in \mathfrak{g}^* \mid \lambda(X_j) = 0, r < j \leq n\}.$$

Clearly $\mathfrak{z}^* \approx \mathfrak{z} \approx \mathbb{R}^r \approx Z$ as topological spaces. Let us also introduce the so-called Pfaffian determinant

$$\begin{aligned} \rho: \mathfrak{z}^* &\rightarrow \mathbb{R}^+ \\ \lambda &\mapsto \left| \det \left(\lambda([X_i, X_j]) \right)_{i,j=r,\dots,n} \right| \end{aligned}$$

and denote with $\mathfrak{g}_j(\lambda)$ the maximal subalgebra \mathfrak{m}_j of \mathfrak{g}_j such that $\lambda([m_j, m_j]) = 0$. Here $[\cdot, \cdot]$ stands for the Lie algebra commutator. With this notation, the following Plancherel theorem holds.

Theorem 6.1. *Let \mathcal{G} be an SI/Z group and let $\Sigma = \{\lambda \in \mathfrak{z}^* \mid \rho(\lambda) > 0\}$. Moreover, let Π_λ be the representation of \mathcal{G} induced by the subalgebra*

$$P(\lambda) = \sum_{j=1}^n \mathfrak{g}_j(\lambda), \quad \lambda \in \Sigma.$$

Then Π_λ is a unitary irreducible representation on $L^2(\mathbb{R}^d)$, and the group Fourier transform

$$\mathcal{F}_\mathcal{G} \varphi(\lambda) = \int_{\mathcal{G}} \varphi(g) \Pi_\lambda(g) dg$$

where dg is the Haar measure of \mathcal{G} , implements an isometric isomorphism - the Plancherel transform -

$$\mathcal{F}_\mathcal{G} : L^2(\mathcal{G}) \rightarrow L^2(\mathfrak{z}^*, \mathcal{HS}(L^2(\mathbb{R}^d)), \rho(\lambda) d\lambda)$$

where $d\lambda$ is the Lebesgue measure on $\mathfrak{z}^ \approx \mathbb{R}^r$.*

This theorem actually provides all we need to define a bracket like (1.6). More precisely we can extend such definition to

$$[\varphi, \psi](\alpha) = \sum_{j \in \mathbb{Z}^r} \langle \mathcal{F}_\mathcal{G} \varphi(\alpha + j), \mathcal{F}_\mathcal{G} \psi(\alpha + j) \rangle_{\mathcal{HS}} \rho(\alpha + j) \quad (6.1)$$

for $\varphi, \psi \in L^2(\mathcal{G})$ and a.e. $\alpha \in \mathbb{T}^r$, the r dimensional torus, that can be identified with $[0, 1)^r$ in \mathbb{R}^r as usual.

Indeed, the bracket (6.1) is related to scalar product of $L^2(\mathcal{G})$ in the same way as (1.6) was, i.e. by means of Plancherel theorem and a periodization over \mathbb{R}^r :

$$\begin{aligned} \langle \varphi, \psi \rangle_{L^2(\mathcal{G})} &= \int_{\mathbb{R}^r} \langle \mathcal{F}_\mathcal{G} \varphi(\lambda), \mathcal{F}_\mathcal{G} \psi(\lambda) \rangle_{\mathcal{HS}(L^2(\mathbb{R}^d))} \rho(\lambda) d\lambda \\ &= \sum_{j \in \mathbb{Z}^r} \int_{\mathbb{T}^r} \langle \mathcal{F}_\mathcal{G} \varphi(\alpha + j), \mathcal{F}_\mathcal{G} \psi(\alpha + j) \rangle_{\mathcal{HS}(L^2(\mathbb{R}^d))} \rho(\alpha + j) d\alpha. \end{aligned}$$

The left translation L is a unitary representation of \mathcal{G} onto $L^2(\mathcal{G})$ and is defined by the group action. The Plancherel transform intertwines the left translation operator L_{xz} , $xz \in \mathcal{G}$ with the representation π_λ : for $\phi \in L^1(\mathcal{G}) \cap L^2(\mathcal{G})$

$$\mathcal{F}_\mathcal{G}(L(xz)\phi)(\lambda) = \pi_\lambda(xz)(\mathcal{F}_\mathcal{G}\phi)(\lambda) = e^{2\pi i \langle \lambda, z \rangle} \pi_\lambda(x) \mathcal{F}_\mathcal{G}\phi(\lambda) |\rho(\lambda)|^{1/2}, \quad \lambda \in \mathfrak{z}^*.$$

By the same arguments of the previous sections, then, one can show that the results of this paper are also true for SI/Z groups when we consider a lattice subgroup Γ of \mathcal{G} given by following form

$$\Gamma = \{\gamma_1 \gamma_2 \in \mathcal{G} : \gamma_1 \in \Gamma_1, \gamma_2 \in \Gamma_2\}$$

where Γ_2 is the integer lattice $\Gamma_2 = \exp \mathbb{Z}X_1 \exp \mathbb{Z}X_2 \cdots \exp \mathbb{Z}X_r$ and Γ_1 is a compatible discrete subset of \mathcal{X} . To avoid repeating the calculations, we leave the proof of the results for the SI/Z groups to the interested readers.

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