

Growth of silicon- and carbon-doped GaAs by chemical beam epitaxy using H₂-diluted DTBSi and CBr₄ precursors

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ABSTRACT

A wide range of n- and p-type doping levels in GaAs layers grown by chemical beam epitaxy is achieved by using H₂-diluted DTBSi and CBr₄ as gas precursors for Si and C, respectively. We show that the doping level can be varied by modifying either the concentration or the flux of the diluted precursor. Specifically, we demonstrate carrier concentrations of 7.8×10^{17} – 1.4×10^{19} cm⁻³ for Si, and 1×10^{17} – 3.8×10^{20} cm⁻³ for C, as determined by Hall effect measurements. The dependence of Si incorporation on the diluted-precursor flux is found to be linear. In contrast, we observe a superlinear behavior for C doping. The dependence of the electron and hole mobility values on the carrier concentration as well as the analysis of the layers by low-temperature (12 K) photoluminescence spectroscopy indicate that the use of H₂ for diluting DTBSi or CBr₄ has no effect on the electrical and optical properties of GaAs.

1. Introduction

Solid-state opto- and electronic devices are typically based on junctions formed by the combination of semiconducting layers with varying electrical conductivities, i. e., with different doping characteristics. The accurate control of the dopant concentration and the possibility of tuning it over a wide range are both essential requirements for the fabrication of certain types of devices, such as heterojunction bipolar transistors and multi-junction solar cells. In the particular case of III-As compound semiconductors, two of the most commonly used elements for n- and p-type doping are Si and C, respectively. These two elements are quite often preferred over other alternatives because of both their comparatively low diffusion coefficients, which facilitate the fabrication of junctions with abrupt doping profiles [1,2], and their high solubility limits [3,4].

In chemical beam epitaxy (CBE), where gas sources - such as metal-organic compounds and hydrides - are used as precursors [5], silane and disilane are the typical gases employed for Si doping. The use of these precursors makes it possible to achieve electron concentrations in GaAs of up to 1 – 5×10^{18} cm⁻³ [6,7]. These values are close to the Si solubility limit at the typical substrate temperatures used in CBE for the growth of this compound (500–600 °C), particularly, about 4×10^{18} cm⁻³ [3,4]. Both

silane and disilane are, however, highly hazardous gases with very high vapor pressures at room temperature. Specifically, they are toxic as well as extremely flammable due to their pyrophoric character. In this context, ditertiarybutylsilane (DTBSi) is an interesting alternative for Si doping since it is a non-toxic and non-pyrophoric liquid that can also be properly refined to meet the purity standards of the semiconductor industry [8,9]. In addition, at room temperature, DTBSi has a subatmospheric vapor pressure, which is high enough for the direct doping of III-V compounds [9]. Surprisingly, even though the first study on the use of this precursor in metal-organic chemical vapor epitaxy (MOVPE) dates from the late nineties [8], we are not aware of any report on the n-type doping of GaAs layers using DTBSi in CBE. With respect to C doping, which has a solubility limit in GaAs grown from the melt of approximately 1.9×10^{19} cm⁻³ [3,4], carbon tetrabromide CBr₄ is one of the common precursors in CBE as well as in gas- and solid-source molecular beam epitaxy (MBE) [10–13]. This precursor provides very high hole concentrations in GaAs, in the range of 3×10^{19} – 5×10^{20} cm⁻³ [10,14–16]. These values actually exceed the solubility limit of C in GaAs, as growth usually proceeds far from thermodynamic equilibrium [4]. Nevertheless, the real challenge when using CBr₄ is to obtain, not high, but low doping levels. So far, two different approaches were reported to achieve moderate C concentrations using this precursor. The most extended one consists in decreasing the CBr₄ flux by

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lowering the temperature of the liquid reservoir, therefore, decreasing its vapor pressure [12]. The second approach is based on controlling the CBr_4 flux by varying the conductance of the gas line in its way from the reservoir to the growth chamber [13]. These approaches yielded hole concentrations down to $1 \times 10^{18} \text{ cm}^{-3}$ [12,13]. However, this value is still too high for the fabrication of certain devices, such as GaAs based solar cells, which demand doping concentrations on the order of $1 \times 10^{17} \text{ cm}^{-3}$ [17,18]. Hence, there is a clear need for the development of more suitable solutions.

In this work, we analyze the possibility of controlling the Si and C doping levels in CBE-grown GaAs layers by using controlled dilutions of DTBSi or CBr_4 in H_2 . H_2 is thus not intended to be used as a simple carrier gas using a bubbler, but to control the actual flux of the dopant precursor. Such an approach is based on previous studies about the growth of GaAs layers doped with Sn [19] and on the use of H_2 -diluted precursors in chemical vapor deposition (CVD) to decrease the risk associated to the use of hazardous gases as well as to improve the performance of amorphous-Si solar cells [20,21]. The analysis of the doping concentration by Hall effect measurements as a function of both the concentration and the flux of the H_2 -diluted precursor demonstrates that, by using this approach, it is possible to modify the Si and C concentrations over at least two and three decades, respectively. Importantly, we also find that the use of H_2 to dilute the precursors has no detrimental effects on either the Hall mobility or the luminescence properties, as shown in the latter case by low-temperature photoluminescence (PL) spectroscopy.

2. Experiment

Epi-ready and semi-insulating GaAs(100) wafers, purchased from Wafer Technology LTD, are used as substrates for the growth of Si- and C-doped GaAs layers. The carrier concentration of the wafers is 7.3×10^6 – $2.8 \times 10^7 \text{ cm}^{-3}$. For our experiments, the as received substrates are In-bonded onto a Mo holder before being loaded into a Riber CBE32 system. A detailed description of the CBE system can be found elsewhere [22]. In order to desorb the protective oxide layer, the substrates are first outgassed at $580 \text{ }^\circ\text{C}$, as measured by an infrared optical pyrometer, in vacuum until observing the characteristic reflection high-energy electron diffraction (RHEED) pattern of GaAs. Subsequently, the substrates are further outgassed for about 5 min at $600 \text{ }^\circ\text{C}$ under the As precursor flux, which results in the observation of a faint (2×4) surface reconstruction. Finally, the substrate temperature is decreased down to $530 \text{ }^\circ\text{C}$ for the growth of either Si- or C-doped GaAs. The growth rate of the GaAs doped layers is set to $0.5 \text{ } \mu\text{m/h}$, as measured by RHEED intensity oscillations frequency in undoped GaAs layers. We note that a certain growth rate reduction could take place during the growth of samples heavily doped with CBr_4 , as this precursor is known to etch GaAs during growth [23,24]. The final thickness of our layers ranges from 0.5 to $1 \text{ } \mu\text{m}$, which was determined assuming the growth rate to be constant. As gaseous precursors for Ga, As, Si and C, we use triethylgallium (TEGa), tertiarybutylarsine (TBAs), DTBSi, and CBr_4 , respectively. Low-temperature ($120 \text{ }^\circ\text{C}$) gas injectors are used for TEGa and CBr_4 , while high-temperature ones are employed for TBAs ($820 \text{ }^\circ\text{C}$) and DTBSi ($1000 \text{ }^\circ\text{C}$) to achieve a high cracking efficiency, as observed by quadrupole mass spectrometry. All gas lines are heated from their reservoir plugs up to the gas injectors to avoid the condensation of the gas precursors along their pathway to the growth chamber. H_2 -diluted precursors are prepared using 99.999% pure H_2 , which is further purified by a heated Ag-Pd membrane. The dilution of DTBSi or CBr_4 is carried out inside a non-heated intermediate reservoir, which is used as dopant source during growth. This reservoir is initially filled with either DTBSi or CBr_4 and then with H_2 until reaching the desired dopant precursor concentration, hereafter referred as dilution factor, which is defined as:

$$f = \frac{P_{\text{dopant}}}{P_{\text{dopant}} + P_{\text{H}_2}}, \quad (1)$$

where P_{H_2} and P_{dopant} are, respectively, the H_2 and the dopant partial pressures inside the reservoir. These partial pressures are measured using a capacitance manometer, whose response is almost independent of the gas type. Precursor fluxes are controlled by a pressure-regulated control system that determines the flux through a hole at the exit of an intermediate chamber located between the precursor container and the growth chamber. The arrival rates of the H_2 -diluted precursors to the substrate surface are given here as beam equivalent pressures (BEP). The BEP values are measured before the growth using an ion gauge, which is placed at the same position as the substrate during the epitaxial process.

Hall effect measurements are performed at room temperature using In contacts and $\approx 1 \text{ cm}^2$ van der Pauw patterns. For the analysis of the Hall data, we take into account the thickness of the surface depletion layer associated to the Fermi level pinning at the GaAs surface. According to Ref. [25], we assume that for n-type GaAs the Fermi level is pinned at 0.68 eV below the conduction band minimum, and for p-type 0.5 eV above the valence band maximum. The analysis of the samples by photoluminescence (PL) spectroscopy is carried out at 12 K using a closed-cycle He cryostat. The PL is excited with an Ar laser, dispersed by a 50-cm Spex monochromator, and detected by a combined Si/(Ga,In)As photodiode. The excitation power density is about 1 kW/cm^2 .

3. Results and discussion

3.1. GaAs doping using H_2 -diluted DTBSi

The electron concentration, n , in GaAs layers doped with H_2 -diluted DTBSi is investigated first as a function of the DTBSi dilution factor, f , as defined by Eq. (1). Fig. 1(a) shows the dependence of the Hall electron concentration on f for a H_2 -diluted DTBSi BEP value of $1.9 \times 10^{-6} \text{ Torr}$. The value of n increases from about 2×10^{18} to $1.4 \times 10^{19} \text{ cm}^{-3}$ as f is varied between 0.03 and 2%. The maximum doping concentration achieved here, $1.4 \times 10^{19} \text{ cm}^{-3}$, is even higher than the highest values reported in the literature for Si doped GaAs layers using DTBSi as gas precursor [8]. When using pure DTBSi (100%), the electron concentration decreases down to $7.4 \times 10^{18} \text{ cm}^{-3}$. We attribute this drop to the self-compensation caused by the formation of $\text{Si}_{\text{Ga}}\text{V}_{\text{Ga}}$ complexes [26] or even Si_{As} due to the amphoteric character of Si in GaAs [1]. Finally, as pointed out in the experimental section, we stress here the necessity of using an elevated temperature for the DTBSi cell in order to achieve a high doping efficiency. Further experiments (not shown here) demonstrate a negligible incorporation of Si as the cell is kept at low temperature ($120 \text{ }^\circ\text{C}$).

Our results demonstrate that the n-type doping level can be modified by varying the dilution factor of the DTBSi gaseous precursor. However, the observed electron concentration spreading data between 0.03 and 1% reveals a poor control on the actual doping level. We attribute the data dispersion to deviations of the actual values of f with respect to the nominal ones. Such deviations are likely to occur due to an uncontrolled partial condensation of DTBSi inside the reservoir during the preparation of the gas mixture. This undesirable effect could be avoided by heating the reservoir above the DTBSi condensation point temperature during the dilution process.

Fig. 1(b) presents the variation of the electron concentration with the BEP of the H_2 -diluted DTBSi precursor for two different values of f , 1 and 2.5%. Since for each series of samples we always use the same gas mixture, these data are not distorted by the possible condensation of DTBSi inside the reservoir during its dilution in H_2 . Regardless of the f value, we observe a monotonic increase of n with the BEP. The results shown in Fig. 1(b) demonstrate that, for a given value of f , the final doping level can be accurately controlled by varying the flux of the H_2 -diluted DTBSi precursor. Specifically, the doping level can be

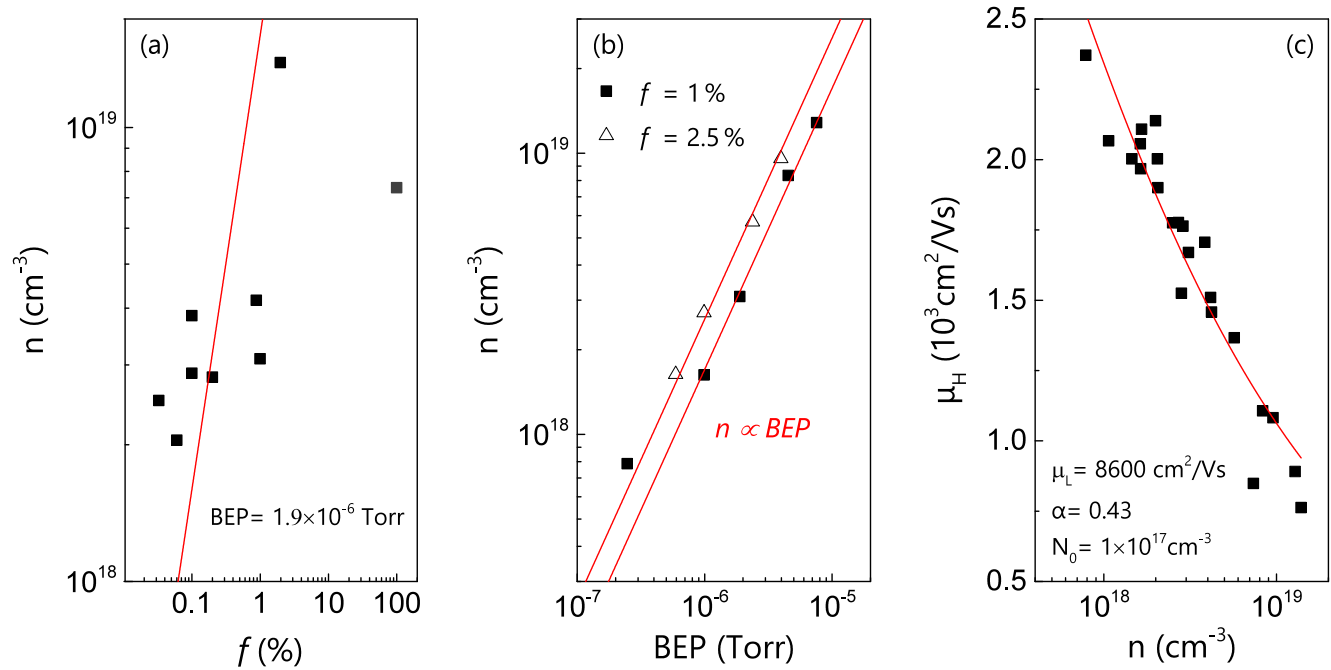


Fig. 1. Electron concentration, as determined by Hall effect measurements, as a function of (a) the dilution factor and (b) the BEP of the H_2 -diluted DTBSi precursor. In (a) the BEP is 1.9×10^{-6} Torr and in (b) the dilution factor is either 1 or 2.5%, as indicated in the legend. The solid lines in (a) and (b) are fits of a first order power law to the experimental data. In (a) the fit is performed for f values lower than 2.5%. (c) Hall electron mobility dependence on the electron concentration. The solid line is a fit of Eq. (2) to the experimental data.

tuned at least from 7.8×10^{17} to $1.3 \times 10^{19} \text{ cm}^{-3}$. Interestingly, for the two different values of f used in this work, the electron concentration increases linearly with the flux of H_2 -diluted DTBSi, as indicated by the corresponding fits shown in Fig. 1(b). Therefore, the incorporation of Si into GaAs as a function of the BEP follows a first order power law. This result reveals that the precursor is cracked in the cell providing active SiH_x radicals to the sample surface. The same linear dependence of n but with respect to the partial pressure of the precursor was reported in Ref. [8] for the MOVPE of GaAs layers doped with pure DTBSi. Such a trend has also been observed in the case of GaN layers doped with DTBSi in MOVPE and CVD [27,28].

Fig. 1(c) shows the dependence of the room-temperature Hall electron mobility on the electron concentration for GaAs doped with H_2 -diluted DTBSi. As the electron concentration increases from 7.8×10^{17} to $1.4 \times 10^{19} \text{ cm}^{-3}$, the mobility decreases from 2.4×10^3 to $0.76 \times 10^3 \text{ cm}^2/\text{V}\cdot\text{s}$. These values are comparable to those reported in Refs. [1,8,29,30] for GaAs layers doped with Si using different precursors including DTBSi. We thus conclude that the dilution of DTBSi in H_2 for the n-type doping of GaAs has no impact on the electron mobility. The dependence of the mobility μ on the electron concentration is well described by the empirical equation proposed by Hilsum [31]:

$$\mu = \frac{\mu_L}{1 + \mu_L/\mu_i} = \frac{\mu_L}{1 + (N/N_0)^\alpha}, \quad (2)$$

where μ_L is the lattice scattering mobility, μ_i the impurity scattering mobility, N the impurity concentration (in this case the concentration of Si, which is assumed to be equal to the electron concentration), N_0 a normalizing impurity concentration, and α a constant with a value lower than unity. As reported in Ref. [31], the normalizing impurity concentration N_0 is 10^{17} cm^{-3} for a wide variety of semiconducting materials, including GaAs. Hence, to fit Eq. (2) to our experimental data, we fix N_0 to 10^{17} cm^{-3} and use μ_L and α as fitting parameters. This fit yields $\mu_L = (8600 \pm 600) \text{ cm}^2/\text{V}\cdot\text{s}$ and $\alpha = 0.43 \pm 0.03$. The obtained value of μ_L is not far from those reported in Refs. [30,31] for Si-doped GaAs layers, namely, $10000 \text{ cm}^2/\text{V}\cdot\text{s}$. Regarding α , the obtained value is close to 0.5, i.e., to the characteristic value found by Hilsum for different n-type

semiconducting materials [31]. Our value of α is furthermore close to 0.4, the value reported in Refs. [30,31] for Si doped GaAs. Based on the above results, we find that Si ionized impurities are responsible for the observed mobility reduction at high electron densities.

3.2. GaAs doping using H_2 -diluted CBr_4

We analyze next the electrical properties of GaAs layers doped with C using H_2 -diluted CBr_4 as gaseous precursor. Fig. 2(a) shows the hole concentration, p , as a function of the CBr_4 dilution factor, f . The BEP of H_2 -diluted CBr_4 is the same for all samples, 6.8×10^{-7} Torr. Despite of the slight spreading of the experimental data, attributed to the partial condensation of CBr_4 inside the reservoir during the preparation of the different gas mixtures (as in the case of DTBSi), we observe a clear superlinear increase of p with increasing values of f . Specifically, p increases from 1×10^{17} to $4 \times 10^{19} \text{ cm}^{-3}$ as f is varied between 0.07 and 1%. The maximum hole concentration achieved here, $3.8 \times 10^{20} \text{ cm}^{-3}$ for a f value of 100%, compares well with the values reported in Refs. [10,14,16] for GaAs layers doped with C using pure CBr_4 as gas precursor. In contrast, to the best of our knowledge, hole concentrations in the 10^{17} cm^{-3} range have never been reported using CBr_4 . The dilution of CBr_4 in H_2 is thus an effective approach for reaching low C concentrations in a controlled fashion.

As demonstrated above for H_2 -diluted DTBSi, the C doping level can also be tuned by varying the BEP of the precursor while keeping constant the CBr_4 dilution factor. Fig. 2(b) shows the variation of p with the BEP of H_2 -diluted CBr_4 for a f value of 0.08%. The hole concentration monotonically increases from $3 \times 10^{16} \text{ cm}^{-3}$ (a value that lays within the range of the typical residual C concentrations obtained in our chamber for non-intentionally doped GaAs layers, from 2 to $4 \times 10^{16} \text{ cm}^{-3}$) to $1.2 \times 10^{19} \text{ cm}^{-3}$ as the BEP is varied between 2.8×10^{-7} and 5.2×10^{-6} Torr. As illustrated by the fit shown in the Fig. 2(b), the dependence of p on the BEP is superlinear and best described by a power law with an exponent close to 4, suggesting the occurrence of complex reactions at the sample surface. Such a power law also properly describes the data presented in Fig. 2(a) for doping

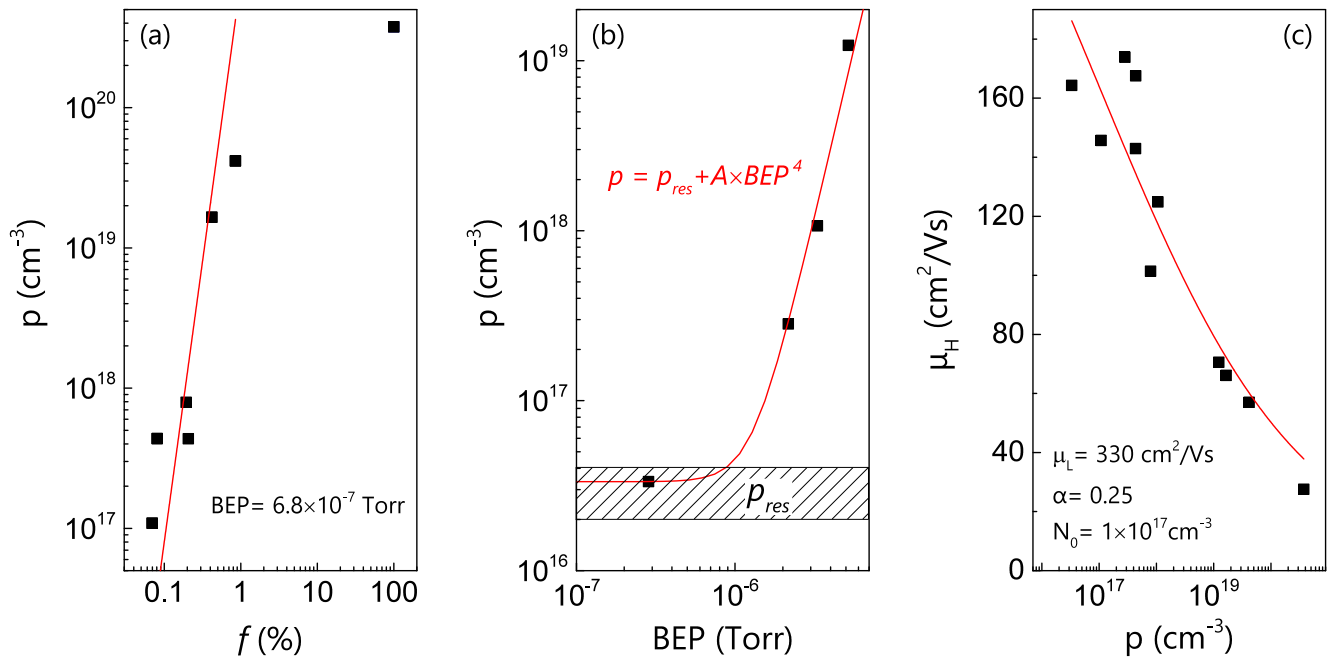


Fig. 2. Hole concentration, as determined by Hall effect measurements, as a function of (a) the dilution factor and (b) the BEP of the H₂-diluted CBr₄ precursor. The BEP in (a) is 6.8×10^{-7} Torr and the dilution factor in (b) is 0.08%. The solid lines in (a) and (b) are fits of a fourth order power law to the experimental data assuming a residual hole concentration p_{res} of 3×10^{16} cm⁻³. The only fitting parameter is thus the power prefactor A . (c) Hall electron mobility dependence on the hole concentration. The solid line is a fit of Eq. (2) to the experimental data.

concentrations below the maximum value. This result is in striking contrast to the approximately linear dependence of C concentration on the CBr₄ supply reported elsewhere for samples grown by metal-organic chemical vapor deposition (MOCVD), MOVPE, MBE and CBE [12,15,24,32,33], with the exception of the work of T. B. Joyce et al. who also observed a superlinear dependence in samples grown by CBE under a high V/III ratio [15]. We tentatively attribute this discrepancy to: (i) the common use of non-identical precursors, noting that CBr₄ could react in concert with, for instance, the Ga precursor or

its byproducts to incorporate C into GaAs [32], and (ii) the intentional supply of H₂ in our work, which is known to at least react with CBr₄ to produce CBr₃ and HBr [24] changing thus the chemistry at the sample surface.

The dependence of the hole mobility on the doping level, as determined by Hall effect, is summarized in Fig. 2(c). The mobility decreases from about 175 to 28 cm²/V·s as the hole concentration increases from 3×10^{16} to 4×10^{20} cm⁻³. The obtained hole-concentration dependent mobilities are comparable to those reported for C-doped GaAs layers

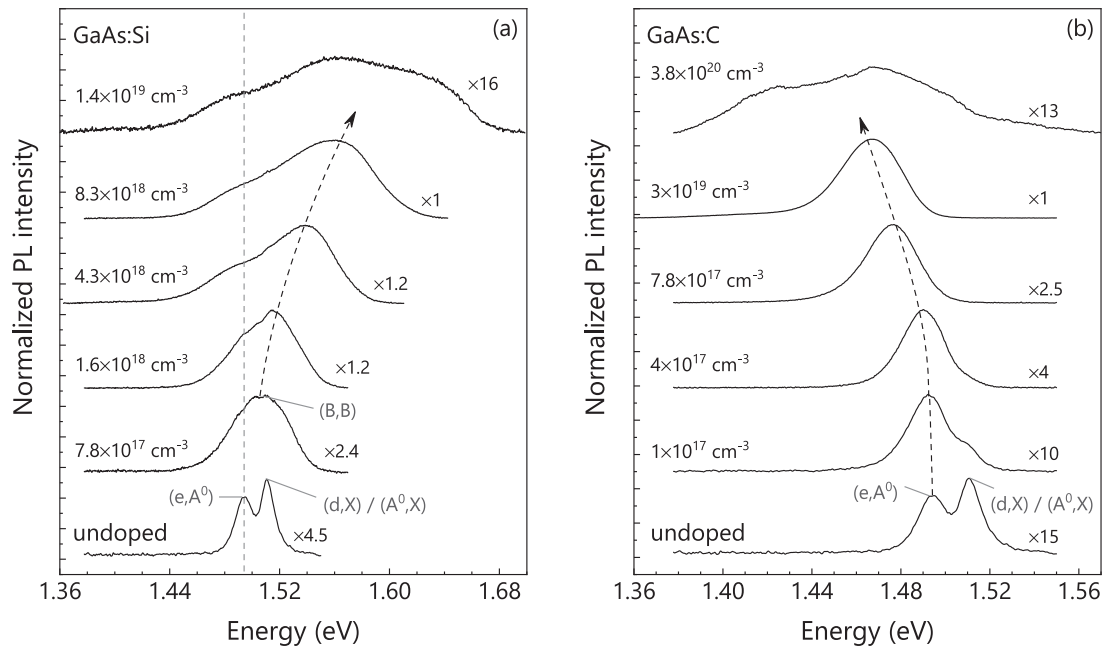


Fig. 3. Low-temperature (12 K) photoluminescence spectra normalized to the maximum intensity and vertically-shifted of GaAs layers doped with different concentrations of (a) Si and (b) C using, respectively, H₂-diluted DTBSi and CBr₄ as gas precursors. In (a) and (b), we also include the spectrum from a non-intentionally doped GaAs layer with a residual C concentration of 3×10^{16} cm⁻³. The normalization intensity factors are indicated at the right side of each spectrum.

grown by either MOVPE or MOMBE using different types of precursors, including the by-products of the chemical reactions occurring between trimethylgallium (TMGa) and AsH₃, CCl₄, or pure CBr₄ [34–37]. As previously obtained for the Si-doped layers, the decreasing mobility with increasing doping level is properly explained considering the scattering caused by ionized C impurities. The best fit of Eq. (2) to our experimental data with N₀ equal to 1 × 10¹⁷ cm⁻³ yields μ_L and α values of (330 ± 20) cm²/V·s and 0.25 ± 0.04, respectively. The values of N₀ and α are similar to those reported in Refs. [35–37] for C-doped GaAs, i. e., N₀ = 3 × 10¹⁷ cm⁻³ and α = 0.3. In view of these results, we conclude that the use of H₂ to dilute CBr₄ has no detrimental effect on the hole mobility, which is exclusively determined by the actual C concentration.

3.3. Luminescence from GaAs layers doped with H₂-diluted DTBSi and CBr₄

We examine the optical properties of the GaAs layers doped with H₂-diluted DTBSi or CBr₄ by low-temperature (12 K) PL. Fig. 3(a) shows the PL spectra of GaAs layers doped with different Si concentrations by using H₂-diluted DTBSi. We also include the PL spectrum of a non-intentionally doped GaAs layer grown in the same system, which exhibits a residual C concentration of 3 × 10¹⁶ cm⁻³. For clarity purposes, the spectra are normalized to the maximum intensity and vertically shifted. The spectrum of the reference sample has two different peaks centered at approximately 1.494 and 1.511 eV. The peak at 1.494 eV is attributed to a band-to-acceptor (e, A⁰) transition, caused by the residual incorporation of C, and the one at 1.511 eV to excitons bound to neutral acceptors (A⁰, X) and/or (d, X) neutral point defects [38–40]. For the Si-doped samples, we observe the (e, A⁰) transition together with a second transition at a higher energies. This second transition, assigned to band-to-band recombination (B, B) [41], blue shifts from 1.51 to 1.56 eV and broadens as the Si concentration increases from 7.8 × 10¹⁷ to 8.3 × 10¹⁸ cm⁻³. For the highest Si concentration analyzed here, 1.4 × 10¹⁹ cm⁻³, the spectrum becomes too broad for a meaningful assessment of the emission energy. Since GaAs becomes degenerated at 12 K for Si concentrations above 2 × 10¹⁸ cm⁻³, the broadening and blue-shift of the (B, B) transition is explained in terms of the Burstein-Moss effect [41]. Such an evolution of the (B, B) transition with increasing Si doping level is characteristic of GaAs:Si layers regardless of the growth technique [41, 42].

Fig. 3(b) presents the normalized to the maximum intensity and vertically-shifted PL spectra of GaAs layers doped with different levels of C using H₂-diluted CBr₄. For comparison, we present here again the spectrum of the non-intentionally doped GaAs layer. In accordance to Refs. [43–45], the PL spectra of the C-doped layers are dominated by the (e, A⁰) transition. Excluding the sample with the highest C concentration (3.8 × 10²⁰ cm⁻³), which exhibits a very broad emission, the (e, A⁰) transition monotonically red shifts and broadens with increasing C concentration. Specifically, the peak energy shifts from 1.494 to 1.467 eV when increasing the C concentration from the residual value up to 3 × 10¹⁹ cm⁻³. The red-shift of the (e, A⁰) transition is due to the band gap narrowing [46, 45, 47], while the increasing linewidth is caused by the Burstein-Moss effect and the relaxation of the k-selection rules [46]. Both the red-shift and the broadening of the peak are typical of highly C doped samples [44, 46, 45, 47]. Therefore, as in the case of the Si-doped GaAs layers doped with H₂-diluted DTBSi, the luminescence properties of the GaAs doped with H₂-diluted CBr are comparable to those of equivalent samples prepared by other techniques using different dopant sources. Consequently, the use of H₂ to dilute either DTBSi or CBr₄ does not have any particular effect on the optical properties of Si and C doped GaAs layers grown by CBE. It is also important to highlight the observed increase on the PL intensity when increasing either the Si or C doping level; this general trend may be explained on the basis of the larger available electron concentration on Si-doped samples, and by an increasing number of recombination centers on C-doped samples.

4. Summary and conclusions

We have demonstrated the n- and p-type doping of GaAs layers in CBE using H₂-diluted DTBSi and CBr₄ as gas precursors. The dilution of pure DTBSi and CBr₄ in H₂ allowed us to achieve a wide range of doping levels. Specifically, GaAs electron and hole concentrations were varied from 7.8 × 10¹⁷ to 1.4 × 10¹⁹ cm⁻³ and from 1 × 10¹⁷ to 3.8 × 10²⁰ cm⁻³, respectively. In the case of H₂-diluted DTBSi, the incorporation of Si into GaAs is found to be linear with respect to the dopant flux. This dependence is analogous to the variation of the Si concentration with the partial pressure of pure DTBSi in MOVPE. In contrast, the dependence of the C incorporation on the flux of uncracked H₂-diluted CBr₄ follows not a linear but a superlinear dependence, revealing the occurrence of complex reactions at the sample surface and making the doping process extremely sensitive to the actual CBr₄ supply. According to our results, for a given dilution factor of either DTBSi or CBr₄ in H₂, it is possible to obtain an accurate control on the doping level by varying the flux of the H₂-diluted precursor. However, we did not achieve a proper degree of control when varying the precursor dilution factor while keeping constant the flux. This result is attributed to the unintentional partial condensation of the precursor inside the reservoir during the dilution process, which results in a poor control on the actual precursor dilution factor. Consequently, a better degree of control is expected by improving the dilution process, such as avoiding condensation by heating the reservoir above the temperature of the pure precursor container when approaching the precursor partial vapor pressure inside the reservoir (as the gas mixture is prepared). Finally, the analysis of the samples by Hall effect measurements and low-temperature PL do not indicate any detrimental or particular effect associated to H₂. Therefore, the use of H₂-diluted precursors allows to tune the doping level over a wide range without affecting the electrical and optical properties of GaAs. We expect that this approach to widen the range of achievable doping levels could be extended to other precursors and material systems.

CRedit authorship contribution statement

K. Ben Saddik: Investigation, Methodology, Conceptualization, Validation, Formal analysis, Writing - Original Draft. **F. Braña:** Investigation, Methodology, Writing - Review & Editing. **N. López:** Investigation, Methodology, Writing - Review & Editing. **J. García:** Investigation, Conceptualization, Methodology, Writing - Review & Editing, Supervision, Project administration, Funding acquisition. **S. Fernández-Garrido:** Investigation, Conceptualization, Methodology, Writing - Original Draft, Supervision.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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References

- [1] N. Furuhashi, K. Kakimoto, M. Yoshida, T. Kamejima, Heavily Si-doped GaAs grown by metalorganic chemical vapor deposition, *J. Appl. Phys.* 64 (9) (1988) 4692.
- [2] B.T. Cunningham, L.J. Guido, J.E. Baker, J.S. Major, N. Holonyak, G.E. Stillman, Carbon diffusion in undoped, n-type, and p-type GaAs, *Appl. Phys. Lett.* 55 (7) (1989) 687.
- [3] O. Madelung, *Semiconductors-Basic Data*, Springer Berlin Heidelberg, Berlin, Heidelberg, 1996.
- [4] R.C. Newman, The upper limits of useful n- and p-type doping in GaAs and AlAs, *Mat. Sci. Eng. B-Adv.* 66 (1–3) (1999) 39.
- [5] W.T. Tsang, Current status review and future prospects of CBE, MOMBE and GSMBE, *J. Cryst. Growth* 107 (1–4) (1991) 960.
- [6] K. Kimura, S. Horiguchi, K. Kamon, M. Shimazu, M. Mashita, M. Mihara, M. Ishii, Silicon doping from disilane in gas source MBE of GaAs, *J. Cryst. Growth* 81 (1–4) (1987) 276–280.
- [7] H. Heinecke, K. Werner, M. Weyers, H. Lüth, P. Balk, Doping of GaAs in metalorganic MBE using gaseous sources, *J. Cryst. Growth* 81 (1–4) (1987) 270–275.
- [8] S. Leu, H. Protzmann, F. Höhnsdorf, W. Stolz, J. Steinkirchner, E. Hufgard, Si-doping of MOVPE grown InP and GaAs by using the liquid Si source ditertiarybutyl silane, *J. Cryst. Growth* 195 (1998) 91.
- [9] Michal Fulem, Květoslav Růžička, Vlastimil Růžička, Tomislav Šimeček, Eduard Hulicius, Jiří Pangrác, Jutta Becker, Jörg Koch, Vapor Pressure of Di-tert-butylsilane, *J. Chem. Eng. Data* 50 (5) (2005) 1613.
- [10] T.J. de Lyon, N.I. Buchan, P.D. Kirchner, J.M. Woodall, G.J. Scilla, F. Cardone, High carbon doping efficiency of bromomethanes in gas source molecular beam epitaxial growth of GaAs, *Appl. Phys. Lett.* 58 (5) (1991) 517.
- [11] K. Zhang, W. Hwang, D.L. Miller, L.W. Kapitan, Carbon doping of GaAs and (In, Ga) As in solid source molecular beam epitaxy using carbon tetrabromide, *Appl. Phys. Lett.* 63 (17) (1993) 2399.
- [12] Q.F. Huang, S.F. Yoon, K.H. Tan, Z.Z. Sun, R. Zhang, J. Jiang, L.H. Lee, Carbon doping in GaAs using carbon tetrabromide in solid source molecular beam epitaxy, *J. Cryst. Growth* 252 (1–3) (2003) 37.
- [13] Y.-C. Chang, Y. Zheng, J.H. English, A.W. Jackson, L.A. Coldren, Wide-dynamic-range, fast-response CBr₄ doping system for molecular beam epitaxy, *J. Vac. Sci. Technol. B* 28 (3) (2010) 310.
- [14] Y.-M. Hoang, Growth of high-quality p-type GaAs epitaxial layers using carbon tetrabromide by gas source molecular-beam epitaxy and molecular-beam epitaxy, *J. Vac. Sci. Technol. B* 11 (3) (1993) 915.
- [15] T.B. Joyce, S.P. Westwater, P.J. Goodhew, R.E. Pritchard, Growth of carbon-doped GaAs, AlGaAs and InGaAs by chemical beam epitaxy and the application of in-situ monitoring, *J. Cryst. Growth* 164 (1–4) (Jul 1996) 371–376.
- [16] S.P. Westwater, T.J. Bullough, The microstructure and thermal stability of CBE grown heavily carbon doped GaAs, *J. Cryst. Growth* 170 (1–4) (1997) 752–756.
- [17] S.R. Kurtz, A.A. Allerman, E.D. Jones, J.M. Gee, J.J. Banas, B.E. Hammons, InGaAsN solar cells with 1.0 eV band gap, lattice matched to GaAs, *Appl. Phys. Lett.* 74 (5) (1999) 729–731.
- [18] D.J. Friedman, J.F. Geisz, S.R. Kurtz, J.M. Olson, 1-eV solar cells with GaInNAs active layer, *J. Cryst. Growth* 195 (1–4) (1998) 409–415.
- [19] C. Garcia Nunez, D. Ghita, B.J. Garcia, Sn doped GaAs by CBE using tetramethyltin, in: *Proceedings of the 8th Spanish Conference on Electron Devices, CDE'2011, IEEE, 2011*, pp. 1–3. <https://doi.org/10.1109/SCED.2011.5744204>.
- [20] D.V. Tsu, B.S. Chao, S.R. Ovshinsky, S. Guha, J. Yang, Effect of hydrogen dilution on the structure of amorphous silicon alloys, *Appl. Phys. Lett.* 71 (10) (1997) 1317–1319.
- [21] Guozhen Yue, Jessica M. Owens, Jennifer Weinberg-Wolf, Daxing Han, Jeffrey Yang, Kenneth Lord, Boajie Yan, Subhendu Guha, Effects of Hydrogen Dilution on a-Si: H and its Solar Cells Studied by Raman and Photoluminescence Spectroscopy, in: *MRS Proc.*, vol. 664, 2001. <https://doi.org/10.1557/PROC-664-A9.7>.
- [22] M. Ait-Lhouss, J.L. Castaño, J. Piqueras, Chemical beam epitaxial growth of GaAs from tertiarybutylarsine and triethylgallium precursors, *Mater. Sci. Eng. B* 28 (1–3) (1994) 155.
- [23] T.B. Joyce, T.J. Bullough, T. Farrell, Optical monitoring of the growth of heavily doped GaAs by chemical beam epitaxy and of the in situ etching of GaAs using CBr₄, *Appl. Phys. Lett.* 65 (17) (1994) 2193–2195.
- [24] Kouta Tateno, Yoshitaka Kohama, Chikara Amano, Carbon doping and etching effects of CBr₄ during metalorganic chemical vapor deposition of GaAs and AlAs, *J. Cryst. Growth* 172 (1–2) (1997) 5–12.
- [25] P.H. Holloway, G.E. McGuire, *Handbook of Compound Semiconductors: Growth, Processing, Characterization, and Devices*, Noyes Publications, 1995.
- [26] A. Mohades-Kassai, Compensation in MBE-grown GaAs doped with silicon and beryllium, in: *ICM 2000. Proceedings of the 12th International Conference on Microelectronics. (IEEE Cat. No.00EX453)*, Univ. Tehran, 2000, vol. 513, pp. 341–344.
- [27] C.J. Deatcher, M.-G. Chaowang Liu, L.M. Smith Cheong, Simon Rushworth, Anna Widdowson, I.M. Watson, Silicon Doping of Gallium Nitride Using Ditertiarybutylsilane, *Chem. Vap. Deposition* 10 (4) (2004) 187–190.
- [28] W.K. Fong, K.K. Leung, C. Surya, Si doping of metal-organic chemical vapor deposition grown gallium nitride using ditertiarybutyl silane metal-organic source, *J. Cryst. Growth* 298 (2007) 239–242, <https://doi.org/10.1016/j.jcrysgro.2006.10.024> (SPEC. ISS).
- [29] M. Druminski, H.-D. Wolf, K.-H. Zschauer, K. Wittmaack, Unexpectedly high energy photoluminescence of highly Si doped GaAs grown by MOVPE, *J. Cryst. Growth* 57 (2) (1982) 318–324.
- [30] M. Kamp, G. Mörsch, J. Gräber, H. Lüth, Te doping of GaAs using diethyl-tellurium, *J. Appl. Phys.* 76 (3) (1994) 1974–1976.
- [31] W.K. Fong, K.K. Leung, C. Surya, Simple empirical relationship between mobility and carrier concentration, *Electron. Lett.* 10 (13) (1974) 259.
- [32] N.I. Buchan, T.F. Kuech, G. Scilla, F. Cardone, Carbon incorporation in metalorganic vapor phase epitaxy grown GaAs using CHyX₄ - y, TMG and AsH₃, *J. Cryst. Growth* 110 (3) (1991) 405–414.
- [33] Jae-Chun Han, Jong-In Song, Sung-Woong Park, Deokha Woo, Growth of ultrahigh carbon-doped InGaAs and its application to InP/InGaAs(C) HBTs, *IEEE Trans. Electron Devices* 49 (1) (2002) 1–6.
- [34] L. Wang, N.M. Haegel, Optical Characterization of Heavily Carbon Doped GaAs, *MRS Proc.* 240 (1991) 87.
- [35] S. Il Kim, Y. Kim, M.S. Lee, M.-S. Kim, S.-K. Min, C. Lee, Hall mobility and temperature dependent photoluminescence of carbon-doped GaAs, *Solid State Commun.* 88 (9) (1993) 743–746.
- [36] E. Richter, P. Kurpas, D. Gutsche, M. Weyers, Carbon doped GaAs grown in low pressure-metalorganic vapor phase epitaxy using carbon tetrabromide, *J. Electron. Mater.* 24 (11) (1995) 1719–1722.
- [37] C.-S. Son, S.-Il Kim, B.-D. Min, Y. Kim, E.K. Kim, S.-K. Min, I.-H. Choi, Electrical Properties of Heavily Carbon-Doped GaAs Epilayers Grown by Atmospheric Pressure Metalorganic Chemical Vapor Deposition Using CBr₄, *Jpn. J. Appl. Phys.* 35 (1996) 6562–6565, <https://doi.org/10.1143/JJAP.35.6562> (Part 1, No. 12B).
- [38] S.-J.P. Park, J.-R.R. Ro, J.-K.S. Sim, E.-H.L. Lee, Growth of GaAs by Chemical Beam Epitaxy Using Unprecracked Arsine and Trimethylgallium, *ETRI J.* 16 (3) (1994) 1–10.
- [39] L. Pavesi, M. Guzzi, Photoluminescence of Al_xGa_{1-x}As alloys, *J. Appl. Phys.* 75 (10) (May 1994) 4779–4842.
- [40] F. Briones, D.M. Collins, Low temperature photoluminescence of lightly Si-doped and undoped MBE GaAs, *J. Electron. Mater.* 11 (4) (1982) 847–866.
- [41] T. Lideikis, G. Treideris, Near band-gap photoluminescence of the MOCVD-grown heavily Si-doped GaAs, *Semicond. Sci. Technol.* 4 (11) (1989) 938–942.
- [42] L. Pavesi, M. Henini, Photoluminescence investigation of Si-doped GaAs grown by molecular beam epitaxy on non-(100) oriented surfaces, *Microelectron. J.* 28 (8–10) (1997) 717–726.
- [43] M. Ozeki, K. Nakai, K. Dazai, O. Ryuzan, Photoluminescence Study of Carbon Doped Gallium Arsenide, *Jpn. J. Appl. Phys.* 13 (7) (1974) 1121–1126.
- [44] H.D. Chen, M.S. Feng, P.A. Chen, K.C. Lin, J.W. Wu, Photoluminescence of Heavily p-Type-Doped GaAs: Temperature and Concentration Dependences, *Jpn. J. Appl. Phys.* 33 (1994) 1920–1927.
- [45] M.S. Feng, C.S. Ares Fang, H.D. Chen, Bandgap shrinkage of degenerate p-type GaAs by photoluminescence spectroscopy, *Mater. Chem. Phys.* 42 (2) (1995) 143–147.
- [46] L. Wang, N.M. Haegel, J.R. Lowney, Band-to-band photoluminescence and luminescence excitation in extremely heavily carbon-doped epitaxial GaAs, *Phys. Rev. B* 49 (16) (1994) 10976–10985.
- [47] J.S. Lee, I. Kim, B.D. Choe, W.G. Jeong, Y.K. Sin, W.S. Min, Luminescence properties of heavily carbon doped GaAs, *J. Appl. Phys.* 79 (12) (1996) 9278–9282.